

Moduli, Supersymmetry,
and
Cosmology

M. Dine
SUSY 2K

I. Cosmology, the cosmological constant
and the anthropic principal

II. Is low energy supersymmetry
a prediction of string theory?

III. Moduli and their
cosmological problems

Really understanding the problems of vacuum selection and moduli in string theory will require understanding the cosmological constant.

Now: data $\Omega_\Lambda \cong 0.7$ [?]

New ideas:

- (1) Holography [somehow]
- (2) Brane World [Somehow]
- (3) Saltatory [Wilczek et al, after Brown and Teitelboim; Abbott]

Data is uncomfortably suggestive of an anthropic explanation.

Weinberg:

A physicist talking about the anthropic principle runs the same risk as a cleric talking about pornography: no matter how much you say you're against it, some people will think you're a little too interested.

A plausible anthropic explanation

requires: (1) a vast no. of metastable states with different values of Λ ;

$$\Delta\Lambda \ll \Lambda_{\text{obs}}$$

(2) An explanation why we can only exist for $\Lambda \leq \Lambda_{\text{obs}}$.

Weinberg: galaxy formation.

$$\Lambda \leq 100 \Lambda_{\text{obs}}$$

Vilenkin:

$$\Lambda \leq 10 \Lambda_{\text{obs}}$$

(1) Possible in string theory?

No reliable demonstration of a

single metastable state [Non-SUSY]

Proposals:

26

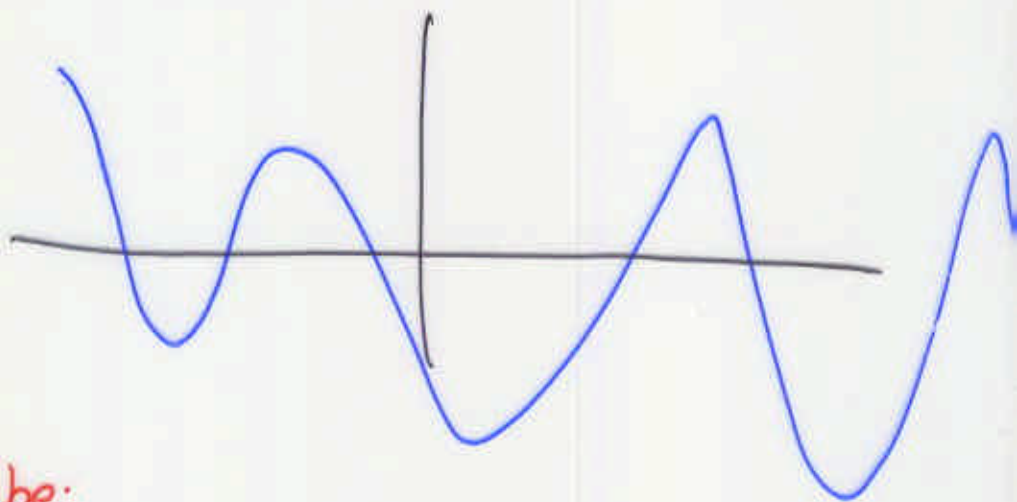
4

"Irrational axion"

$$V = \mu_1^4 \cos(a/f_1) + \mu_2^4 \cos(a/f_2) + V_0$$

f_1/f_2 : irrational; $\mu_1, \mu_2 > V_0$

T. Banks,
(N. Seiberg
M.D.)



Easy to describe;
No known string examples. At best,
several groups, f_i/f_j large primes.

E

...

More promising: 4-form flux

(Bousso - Polchinski
Donoghue)

Quantized: $F_{\mu\nu\rho\sigma}^{[i]} = g^{(i)} n^{(i)} \epsilon_{\mu\nu\rho\sigma}$

[Dirac-Teitelboim;
one for each
3-cycle]

$$E = \sum n_i^2 g^{(i)^2} - \Lambda_0$$

If ~ 120 fluxes, and energy spacing

$\sim \frac{1}{30} M_{Pl}$, dense set of states

with splittings $\ll \Lambda_{obs}$

(no. of states: $\int_{n_i}^{n_i + \Delta n_i} dn n^{N-1} \Omega_N$)

Examining quantization conditions,

$$g_i^2 \sim \frac{1}{R^2} \quad (\text{say})$$

so small g_i from large R .



In this framework, anthropic explanation of all quantities inevitable.

E.g.

$$F_{\mu\nu\rho\sigma}^2 |H|^2 \Rightarrow m_H^2$$

$$F_{\mu\nu\rho\sigma}^2 F_{\mu\nu}^{(i)2} \Rightarrow g^{(i)-2} \quad \text{etc.}$$

(Banks, Moore, Motl, M. D.)

No empty universe problem as in Abbot:

E.g. in last jump, big shift in minimum of inflaton field ($F^2 I^2$ terms in V)

More appealing: very large dimensions.

Topological charges responsible for

stabilizing large dimensions

(as in Arkani-Hamed, Dimopoulos, Dvali, March Russell)

Cancellation of cosmo. const.

requires $d=2$, bulk SUSY

S.M. parameters not very sensitive

to fluxes. Main difficulty:

empty universe. Solvable?

Still, a potentially natural, plausible
use of anthropic principle.

In conclusion, a much more restrained
role conjectured for anthropic principle.

Supersymmetry and String Theory?

Is supersymmetry [softly broken $N=1$] a prediction of string theory?

Last two years: plausible alternative solutions to hierarchy problem.

The more supersymmetry, the more we know: [e.g. that vacua exist]

- | | |
|------------------------------|---|
| (a) $N \geq 1$ in $D \geq 5$ | } Exact moduli spaces exist non-perturbatively [Support from dualities] |
| (b) $N \geq 2$ in $D=4$ | |

(c) $N=1$ in $D=4$?

- Exact moduli spaces exist. Dualities, etc.
- Generically, approximate moduli spaces. Dual descriptions agree where appropriate \Rightarrow Exist

Exact moduli: guaranteed, in some cases, by discrete symmetries, trivial in dynamics. Can often study dual descriptions.

Approximate moduli: generically, superpotentials lift moduli (e.g. $W(s) = e^{-s/b_0}$ for gluino condensation.)

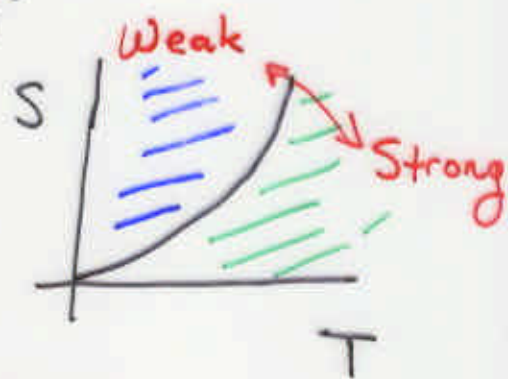
At weak coupling, [Heterotic]

$$S = g^{-2} V / l_s^6 \quad T = R^2 / l_s^2$$

Strong coupling (Horava-Witten)

$$T = V^{1/3} R'' / l_{11}^2 \quad S = V / l_{11}^6$$

One can pass from weak to strong coupling keeping S , T large. Holomorphic quantities agree.



$$f^a = S - n^a T + \mathcal{O}(e^{-S}, e^{-T})$$

$$W = e^{-3S/b_0} [1 + \mathcal{O}(e^{-nS})]$$

[see Kaplunovsky

Waldron
this meeting

Assuming string theory describes nature, it seems likely that the moduli are stabilized on the approximate moduli space (small gauge couplings, hierarchy)

This seems unlikely to be generic.

Proposals:

- Kahler stabilization
- Racetrack (requires particular gauge groups; not generic)



[Cosmological constant]

Predictive power: limited in both cases (holomorphic) but any prediction would be spectacular!

$N=0$

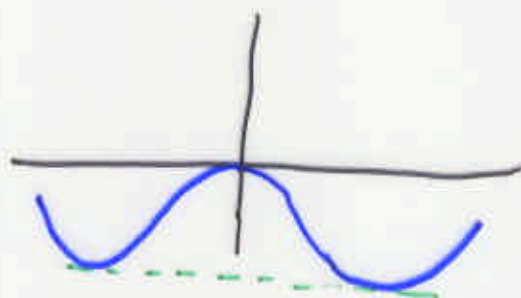
(11)

What distinguishes $N=1, N=0$?

Approximate moduli space and its properties. Asymptotically regions with $4D, N=1$, or not.

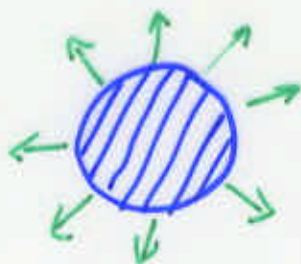
Generically, $N=0$ theories have problems; perhaps the germ of a string prediction?

- Tachyons

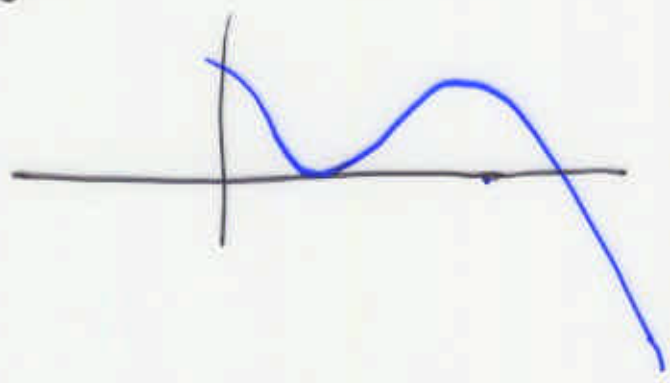


$$V_0 \propto -\frac{1}{g^2} \rightarrow -\infty$$

- Catastrophic vacuum decay
(Witten; Fabinger-Horava)



With our current state of knowlege,
 not decisive. E.g. stabilization in a
 regime with no moduli? Very metastable?



$$\Gamma \sim e^{-(RM)^a}$$

Other (dark horse) possibilities?

Non-perturbative anomalies?

Witten (1985-6): none in field theory
 limit under weak conditions? *

Inherently stringy:

Look for discrete gauge anomalies,
 SU(2) anomalies in a variety of
 theories

* Berkooz, Leigh, Polchinski, Schwarz, Seiberg, Witten:
 examples in open strings ("dual" to
 closed.)

Can we find non-perturbative consistency conditions in closed string theories?

$SU(2)$ (global) anomalies?

Anomalies in discrete symmetries?

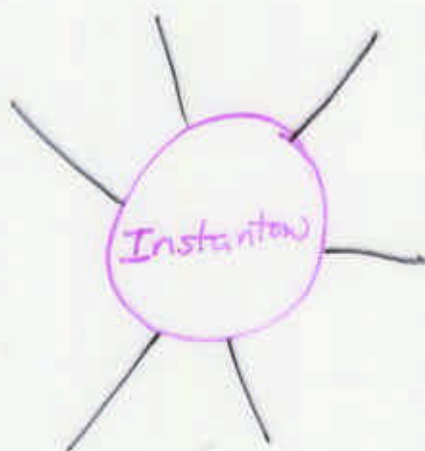
(T. Banks
& M. D.;
D. MacIntyre
& M. D.)

Belief: discrete symmetries in string theory

are gauged, so anomalies would signal an inconsistency.

(Krauss-Wilczek)
Ibanez, Ross; Banks, M. D.

$$\psi\psi\dots\psi \rightarrow \alpha^r \psi\psi\dots\psi$$



Must allow

GS mechanism:

$$a \int F \tilde{F}$$

$$a \rightarrow a + k \frac{2\pi}{N}$$

Discrete symmetries: are they gauged?

36

(74)

(1) Torus



Symmetries $x \rightarrow e^{2\pi i/6} x$

- remnant of general coord. invariance.

(2) $E_8 \times E_8$ $E_8 \leftrightarrow E_8$

subgroup of a continuous group unbroken
on subspace of moduli space.

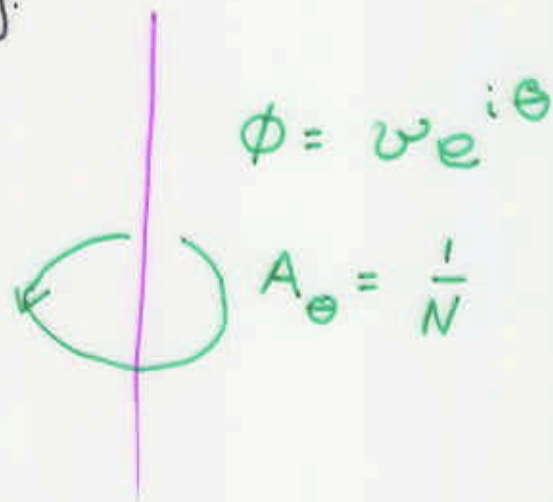
(3) T duality in heterotic string
(as above)

(4) CP in heterotic string

Others not so obvious. Discrete symmetries
of orbifold models? S-dualities?

Another approach: construct cosmic string solutions, ask what happens when bring particle around string.

$$\psi \rightarrow \alpha \psi$$



$$L_2 = m - \frac{2\pi}{N}$$

Analogs in string theory?

"Orbifold cosmic string"

(Bagger, Callan, Harvey)

"Internal" $Z_N, N=2,3,4,6$

Compactify x_2, x_3 on torus with

Z_N symmetry, mod out by

product



Difficulties; alternatives (Banks, Motl, MD)

Previous searches (Banks, M.D.; MacIntyre)
limited. Currently attempting a broader
search (non-SUSY, asymm. orbifolds, ...
(with K. Choi, M. Graesser,
J. Gray)

I.F. Find examples:

Inconsistency?

Evidence some discrete symmetries
not gauged?

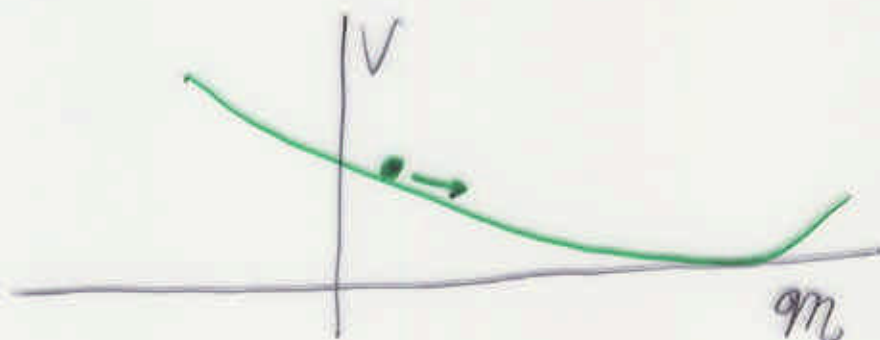
We have hardly proven supersymmetry,
but in spirit of conference, assume
for rest of discussion*, and consider

Cosmology

Moduli + Cosmology

Potential inflatons (Binetruy, Gaillard)
(recently: Banks)

Potential problems: if associated with supersymmetry breaking, tend to store too much energy.



Solutions: Moduli heavy (10^2 of TeV)

[Tuning?] [Baryons - via A.D. (Yanagida et al)]

Late inflation, Enhanced symmetries?

Enhanced symmetries

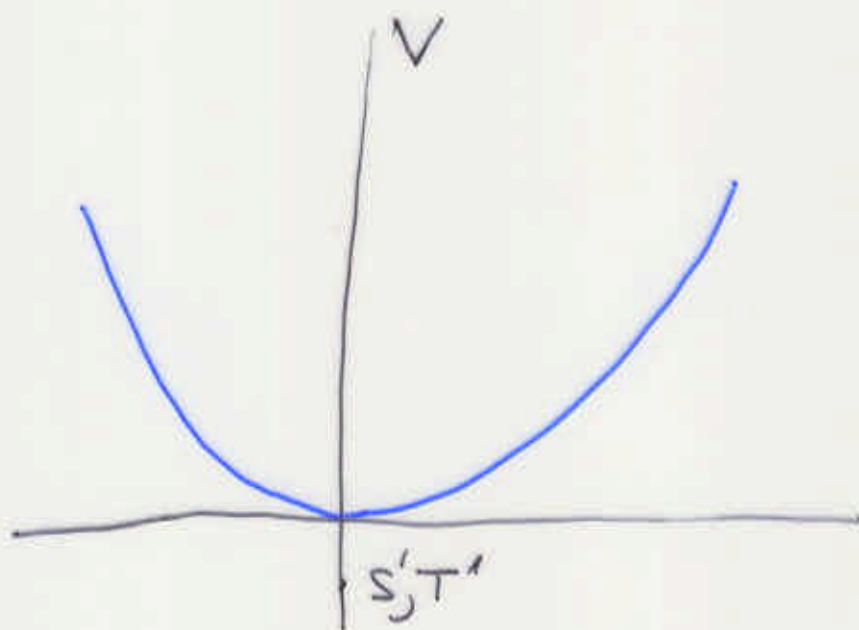
Familiar from studies of duality.

Indeed, pts. in moduli space exist where all moduli transform under broken symmetries. Naturally minima of potential both now + in early universe.

E.g. II_B

Fixed pt. of

$$\tau \rightarrow -1/\tau$$



This provides a natural soln.

(Randall,
Thomas,
M.D.)

to the moduli problem since

enhanced symmetry points are stationary

pts. of effective action (necessarily!)

Candidate minima.

[Giudice, Tkachev, Riotto,
Felder, Lind, Kofman]

But troubling: expect (compare
self-dual pt. of $\vec{E} \leftrightarrow -\vec{B}$ duality)

$\alpha \sim 1$.

Perhaps sometimes simply small

but racetrack models

(Shadmi,
Shirman, Nir,
M.D.)

(esp. Iizawa-Yanagida) suggests

alternative: S fixed with unbroken

SUSY at a high scale. No associated

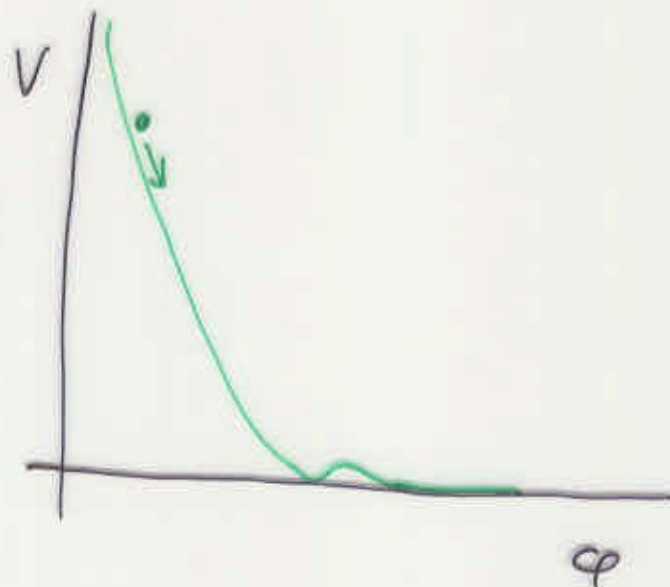
moduli problem.

Suppose stabilization as in racetrack or Kahler:
Brustein - Steinhardt

$$S = e^{\varphi}$$

$$V = e^{-\frac{1}{\alpha} e^{\varphi}}$$

$$\ddot{\varphi} + 3H\dot{\varphi} + V'(\varphi) = 0$$



Inevitably overshoot.

Horne, Moore: moduli space has finite volume.

Finite (not too small) probability that start not too far from minimum.

Banks, Berkooz, Moore, Shenker, Steinhardt:

Kahler potential corrections?

Focus on zero modes ($\bar{k}=0$) inconsistent.

Kahler potential corrections:

Don't help Even allowing arbitrary k_j
 a large no. of tunings. (Nir, Shirman, Shadmi,
 M.D.)

Non-zero momentum?

If zero modes dominate, $p = 0$

$$R \sim t^{1/2} \quad H \sim \frac{1}{3t}$$

$$\ddot{\varphi} + \frac{1}{t} \dot{\varphi} = 0 \quad (\text{neglecting potential})$$

$$\varphi = a \ln t + b \quad \dot{\varphi}^2 \sim \frac{1}{t^2} \sim \frac{1}{R^6}$$

$$(V \sim e^{-t^a} !)$$

Non-zero modes:

$$\ddot{\varphi} + \frac{1}{t} \dot{\varphi} + \frac{k^2}{t^{4/3}} \varphi = 0$$

$$\varphi \sim t^{-1/3} \cos\left(\frac{3}{2} t^{2/3} k\right)$$

$$\dot{\varphi}^2 \sim t^{-4/3} \sim \frac{1}{R^4}!$$

(22)

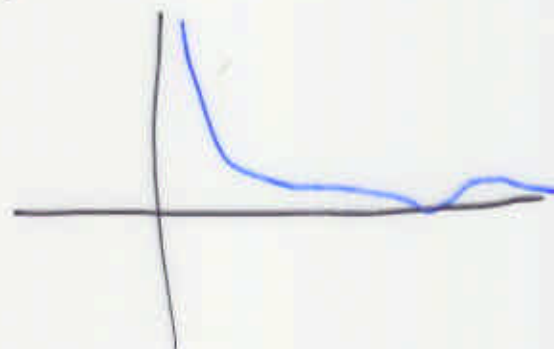
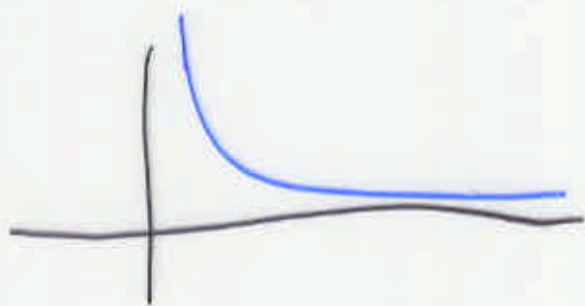
Model for ^{non} zero modes: temperature. (Hug, Steinhardt, Waldra, M.O.)

$$\Omega(g, T) = - \frac{(N^2 - 1)}{24} \pi^4 T^4 \left(1 - \frac{3}{8} \frac{g^2 N}{\pi^2} + \frac{g^3 N^{3/2}}{\sqrt{2}} \pi^3 + \dots \right)$$

For slowly varying g , a potential

for $g = e^\phi$. Valid for $\Lambda(g) \ll T$

$$\Lambda(g) \gg T: \exp\left(-\frac{8\pi^4}{g^2 N}\right)$$



$$p = \frac{1}{3} \rho \quad R^{-1} = \frac{1}{\sqrt{t}} \quad H = \frac{1}{2t}$$

$$\phi \sim \frac{1}{\sqrt{t}} + C \quad \dot{\phi}^2 \sim \frac{1}{t^3}$$

So competition between potential, damping term.

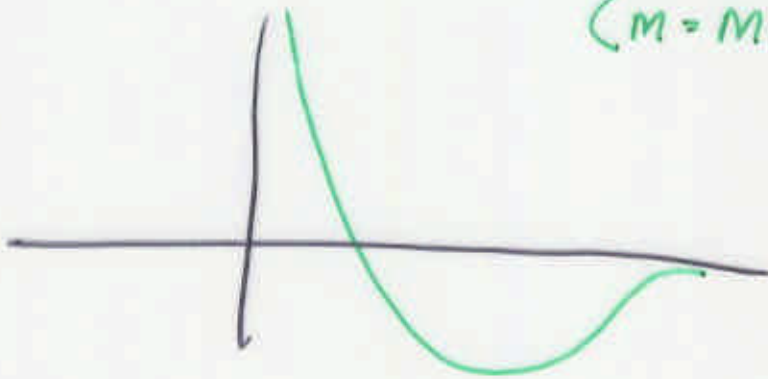
For large N , reasonable range of initial cond:

- soft landing in min.

System need not really be in
thermal equilibrium.

$$\text{IF } V = -g^2 M^4 + V_{np}$$

($M = M(t)$)



- system easily finds true minimum.

So a picture for modulus stabilization
+ solution of the moduli problems
of string cosmology.

But many questions:

- Cosmological constant
- Inflation
- Vacuum selection
(why not $N > 1$, $d > 4$?)

⋮

Conclusions: Anthropic "Lite"

Suppose a smaller set of vacuum states, $N > 1$, $d > 4$, $N = 1$ (approx. moduli), etc. "Early" universe samples all.

Perhaps generic states can't grow large, develop structure.

E.g. $N \geq 2$ exact moduli space.

Non-zero modes dominate evolution.

(Always radiation dominated)

$N = 1$: metastable vacua not generic [?].

⋮

If we understood Λ , moduli stabilization, rest might not be so hard.