

Leading Baryon Production in DIS and Photoproduction

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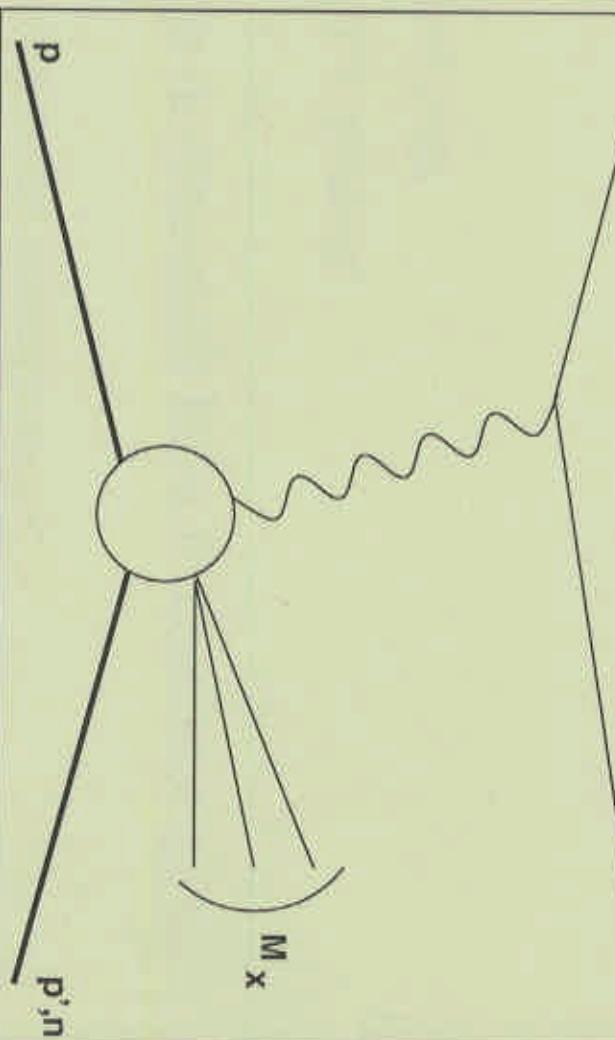
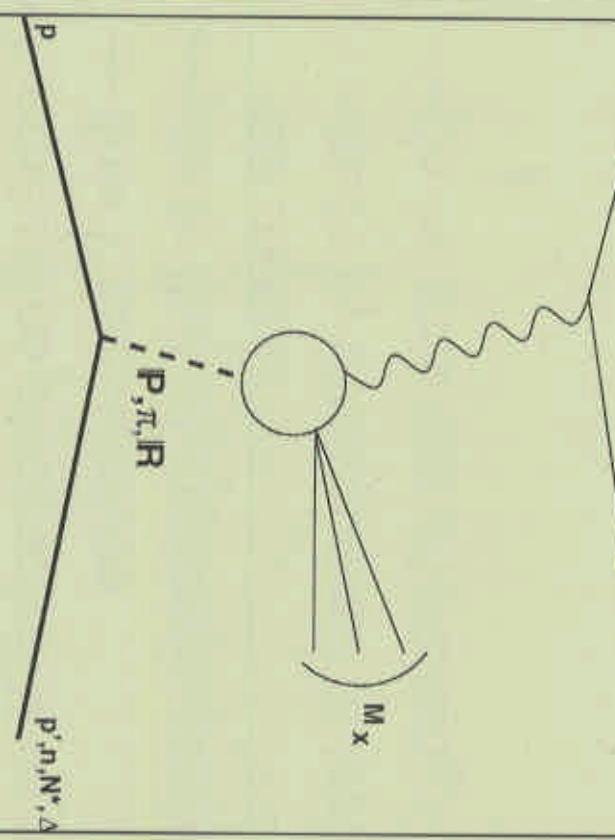
Overview

- Introduction
- Motivation
- Detectors
- Individual Analyses
 - Leading Neutron: neutron energy spectra (Abstract 883)
 - Leading Neutron: dijet photoproduction (Abstracts 967 and -)
 - Leading Neutron in DIS: neutron p_T distributions (Abstract 882)
 - Leading Proton (Abstract 962)

Introduction

hadronization of proton remnant

$e \rightarrow e' + M_X$



Kinematic Quantities

$$x_L = z = \frac{E_{LB}}{E_p}$$

- factorization:
 $\sigma(ep \rightarrow e'nX) = f_{\pi/p} \sigma(e\pi \rightarrow e'X)$
- rescattering effects

$$t = (p_{LB} - p_p)^2 = -\frac{p_T^2}{x_L} - \frac{(1 - x_L) \cdot (m_n^2 - m_p^2 x_L)}{x_L}$$

Motivation

- study production mechanisms for leading baryons
- if production mechanism is an exchange mechanism:
 - check factorization
 - study flux of exchange particle in incoming proton, compare different models
 - study p_T -distributions of leading baryon
 - probe structure of exchange particle

Detectors

- leading protons:
 - employ standard techniques: Roman pots
 - H1 FPS: 2 stations, $\theta_p < 0.5$ mrad
 - ZEUS LPS: 6 stations
- leading neutrons: angular acceptance: $\theta_n < 0.8$ mrad
 - H1: Forward Neutron Calorimeter FNC
 - ZEUS: FNC + Forward Neutron Tracker FNT

Leading Neutron Production

ZEUS Preliminary

- fractional neutron yield

increases with Q^2 , saturates

$\Delta Q^2 < 0.02 \text{ (GeV}^2)$

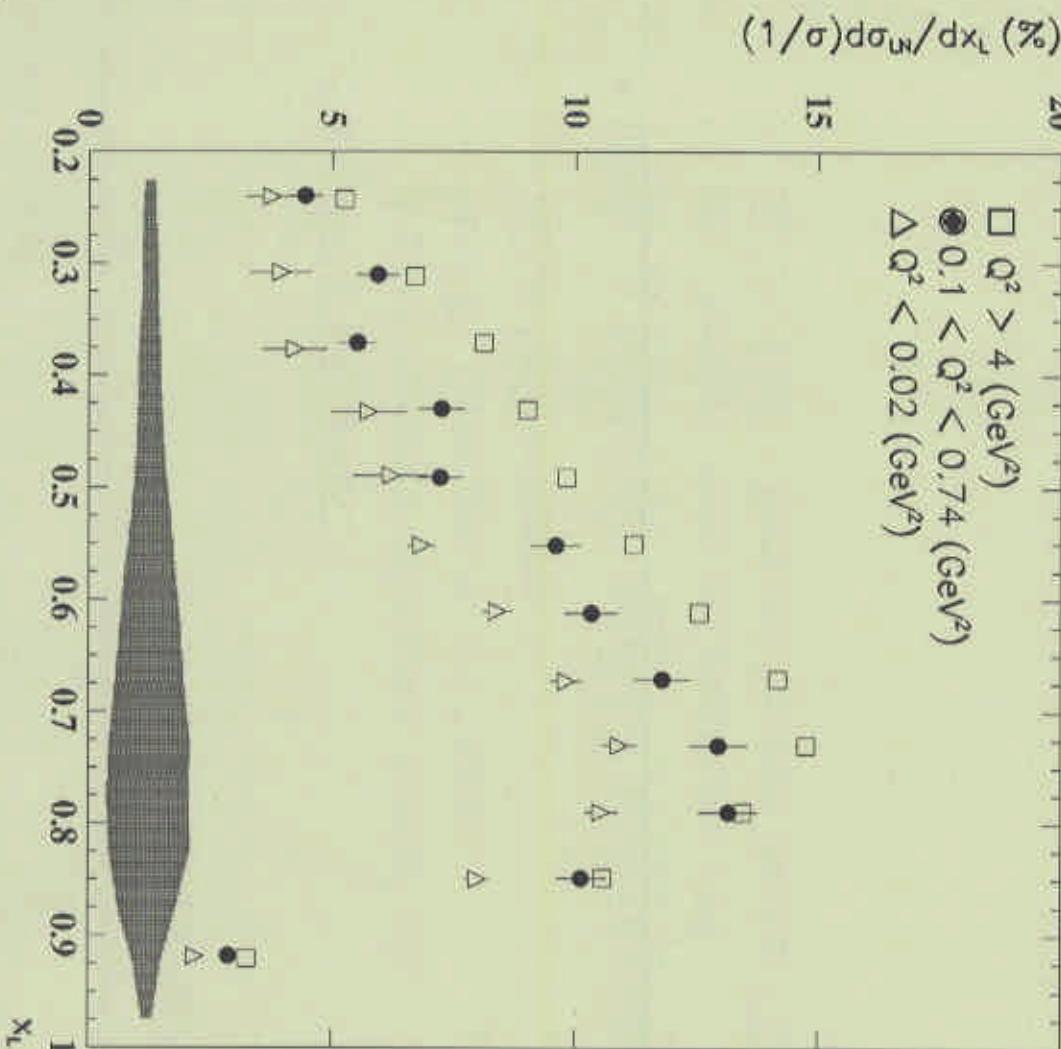
at about 4 GeV^2 (not shown)

- can be attributed to absorp-

tive effects: as Q^2 increases,
size of virtual photon

decreases, probability of
rescattering resulting in loss
of neutron decreases

- shaded band shows system-
atic uncertainty due to FNC
acceptance uncertainty



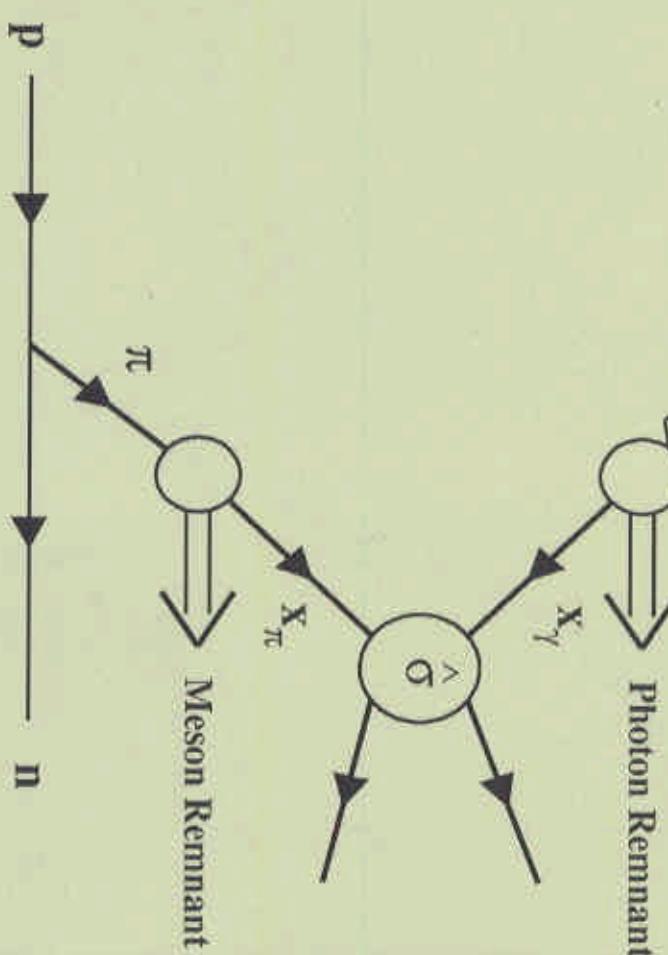
Dijet Photoproduction with LN

$e \rightarrow \gamma \rightarrow \text{jet jet } X$

- H1:
 - tagged photoproduction
 - Jets: cone algorithm with $R=1$, $E_T > 4(7)$ GeV, $-1 < \eta < 2$
 - neutron: $E_n > 400$ GeV, $\theta_n < 0.8$ mrad

- ZEUS:

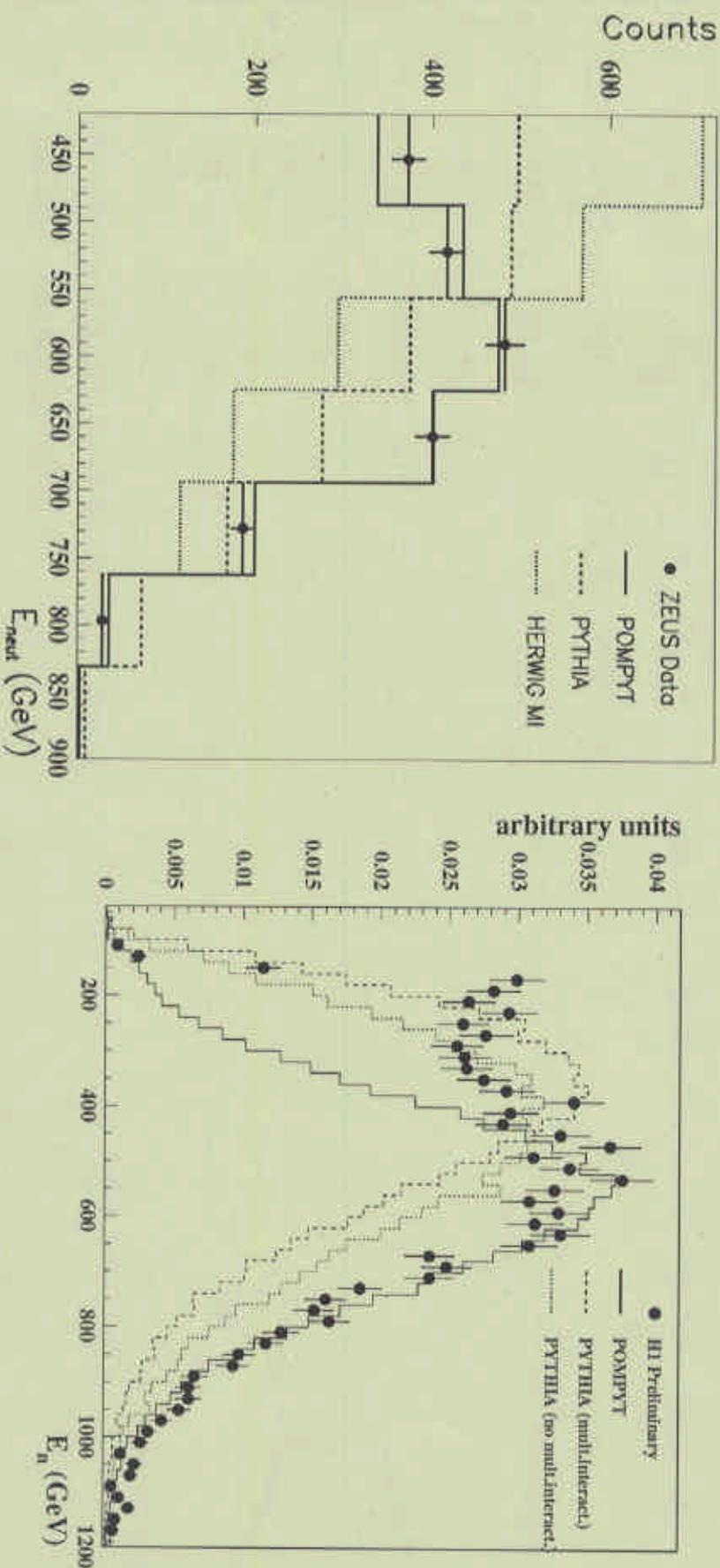
- $0.2 < y < 0.8$
- Jets: k_T algorithm, $E_T > 6$ GeV, $-2 < \eta < 2$
- neutron: $E_n > 400$ GeV, $\theta_n < 0.8$ mrad
- measured: neutron energy spectra, differential dijet cross section as function of E_T , η , x_π
- goals: n production mechanism, check factorization, compare pion flux models



Abstracts 967 and -

Leading Neutron Production Mechanisms

ZEUS 1995 Preliminary



for $x_L > 0.5$ MCs with one-pion-exchange (POMPYT, RAPGAP) describe shape (and normalization) of neutron energy spectrum, inclusive MCs do not

Dijet Photoproduction with LN (*cont'd*)



$$OBS = \frac{\sum_{jet} E_T e^\eta}{2E_\pi}$$

OPE also describes the jet energy spectrum independent of



factorization: shape of neutron energy spectrum independent of

$$x_\gamma^{\text{OBS}} = \frac{\sum_{jet} E_T e^{-\eta}}{2E_\gamma}$$

A plot showing x_γ^{OBS} versus E_n (GeV). The x-axis ranges from 400 to 900, and the y-axis ranges from 0 to 100. Data points are shown for different ranges of x_γ^{obs} : $0.1 < x_\gamma^{\text{obs}} < 0.4$ (solid circles), $0.4 < x_\gamma^{\text{obs}} < 0.8$ (open diamonds), $0.8 < x_\gamma^{\text{obs}} < 1.0$ (open squares).

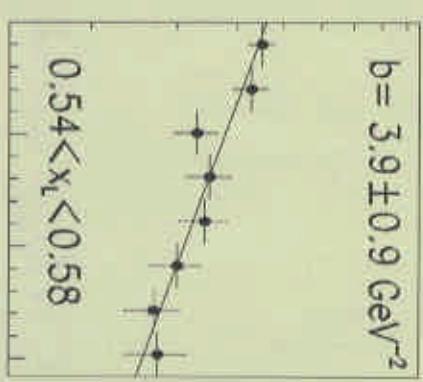
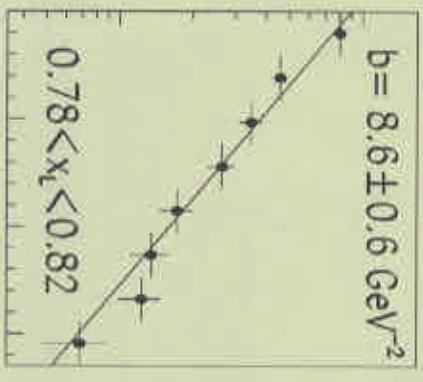
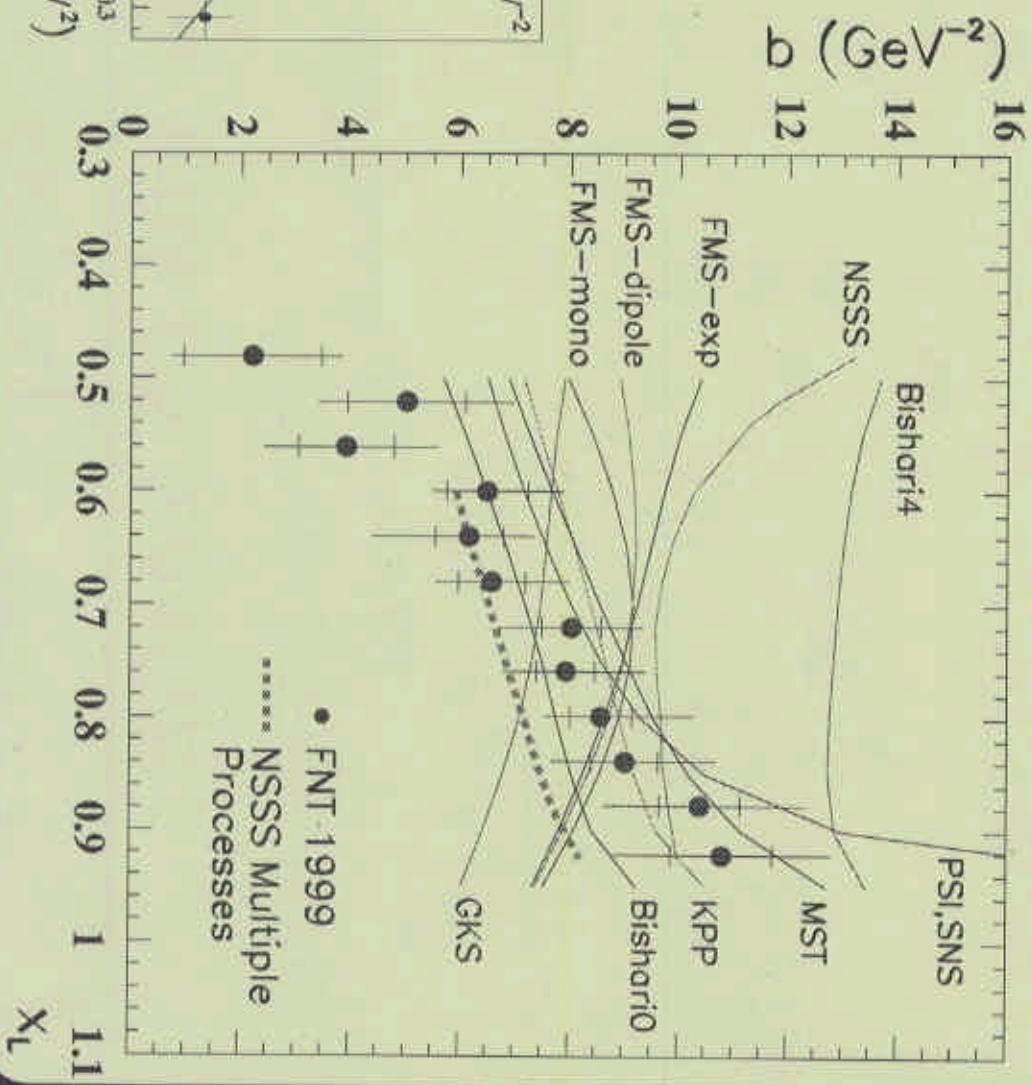
Abstracts 967 and -

Leading Neutron p_T Distributions in DIS ZEUS PRELIMINARY 1999

$$\frac{dN}{dp_T^2} \propto e^{-b(x_L)p_T^2}$$

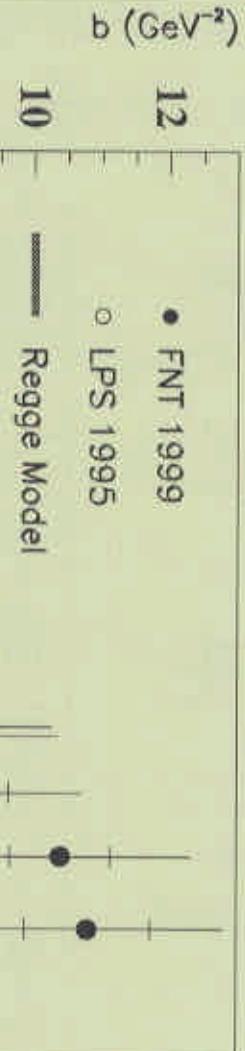
- measurement uses the ZEUS Forward Neutron Tracker (FNT)
- slope b is a function of x_L

ZEUS PRELIMINARY 1999

 dN/dp_T^2 (arb. units) $b = 3.9 \pm 0.9 \text{ GeV}^{-2}$  $b = 8.6 \pm 0.6 \text{ GeV}^{-2}$ 

Abstract 882

Leading Neutron p_T Distributions in DIS ZEUS PRELIMINARY



- comparison with leading proton p_T distributions
- for leading protons, iso-scalar exchanges can also contribute

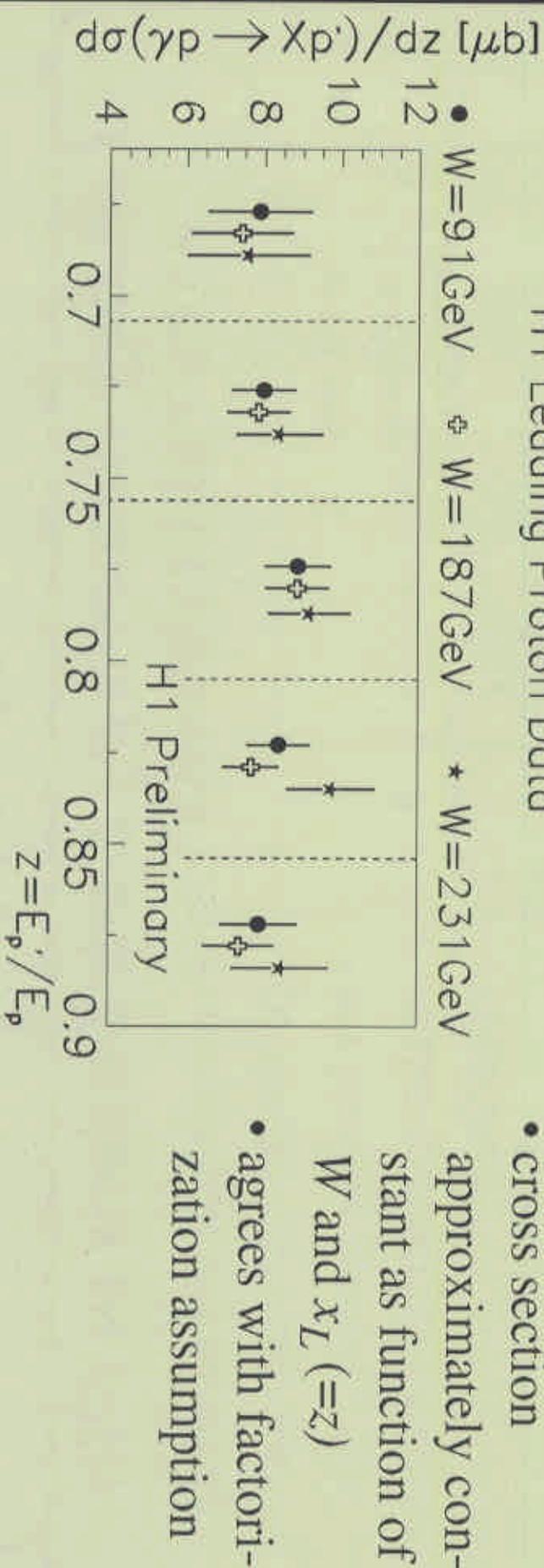
Leading Proton in Photoproduction

- using H1 Forward Proton Spectrometer (FPS), tagged photoproduction, 1996 data

- $p_T < 0.2$ GeV

- $0.66 < x_L < 0.90$: in this region diffractive process suppressed relative to π , reggeon exchange

H1 Leading Proton Data



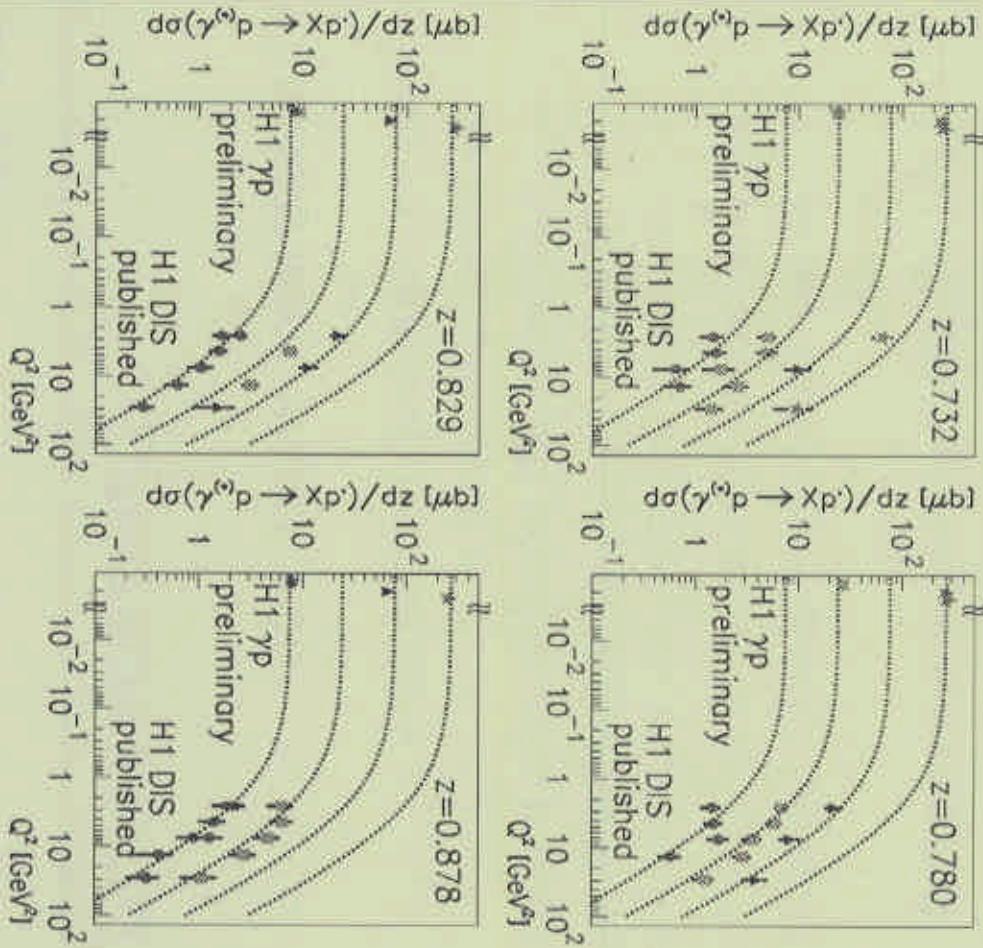
- cross section

approximately constant as function of W and x_L ($= z$)

- agrees with factorization assumption

Leading Proton *cont'd*

H1 Leading Proton Data
 • $M_x < 40\text{GeV}$ ($\times 1$)
 ■ $40\text{GeV} < M_x < 60\text{GeV}$ ($\times 3$)
 * $60\text{GeV} < M_x < 80\text{GeV}$ ($\times 9$)
 * $80\text{GeV} < M_x < 100\text{GeV}$ ($\times 30$)



- model F_2 as sum of VMD term and partonic term with 2 free parameters, Q_{VM}^2 , C_{VM}

$$F_2^{LP} = F \left(C_{VM}^{LP} C_{VM} \cdot F_2^{VMD} + \frac{Q^2}{Q^2 + Q_{VM}^2} \cdot F_2^{part} \right)$$

- use Q_{VM}^2 , C_{VM} from fit to inclusive F_2 , use global scale factor F

- fit with C_{VM}^{LP} as free parameter: $C_{VM}^{LP} = 0.23 \Rightarrow$
 - VMD part suppressed in LP events compared to inclusive

Summary and Conclusions

- Factorization works well
- Leading neutron production:
 - in dijet photoproduction
 - MCs with OPE model describe neutron spectrum and jet distributions for $x_L > 0.5$
 - diffraction plays small role in that domain
 - data sensitive to pion flux, insensitive to pion PDF
 - p_T^2 distributions well described by $\exp(-b(x_L) p_T^2)$
 - slope b rises linearly with x_L ; measured dependence compared to various exchange models
 - yield of leading neutrons as function of Q^2
 - at low Q^2 quick rise, saturation above 4 GeV²
- Leading proton production
 - cross section approximately constant as function of W and x_L
 - VMD part of LP cross section appears suppressed compared to inclusive cross section