

Dynamical gluon mass corrections

in heavy quarkonia decays

A. A. Natale - *IFT - Unesp*

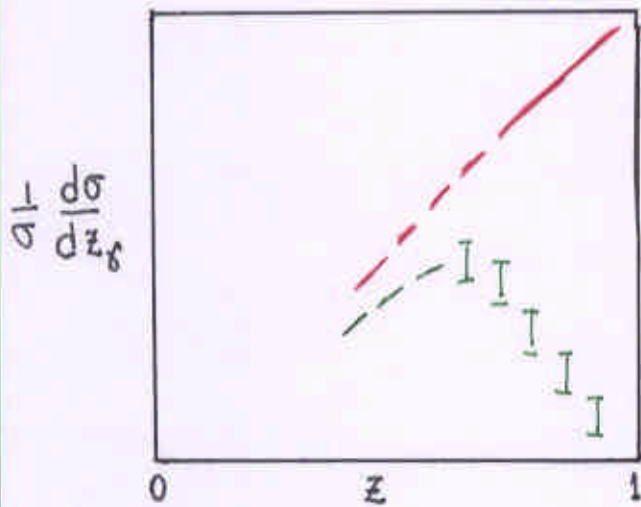
hadronic/radiative decays of quarkonia

- $\Gamma(V \rightarrow \text{hadrons})$
- $\Gamma(V \rightarrow \gamma + X)$, $V = \Upsilon, J/\Psi$
- measure α_s and test perturbative QCD

(a) consistency of α_s measurements at m_Υ and $m_{J/\Psi}$

(b) understand photon spectrum in $V \rightarrow \gamma + X$

photon spectrum at $z(\equiv 2E_\gamma/W) \rightarrow 1$ (signal of nonperturbative QCD)



→ perturbative QCD

→ nonperturbative QCD

One possible explanation: parametrize nonperturbative contributions in terms of an effective gluon mass (M_g)

Parisi and Petronzio, PLB94(80)51; Consoli and Field, PRD49(94)1293; Field, NPB54A(97)247, JPG23(97)41

optimal values:

$J/\Psi \rightarrow M_g \sim 0.66 \text{ GeV}$; $\Upsilon \rightarrow M_g \sim 1.2 \text{ GeV}$

Dynamical gluon mass?

Nonperturbative solution of Schwinger-Dyson equations for the gluon propagator (Cornwall, PRD26 (82) 1453)

$$D(P_E^2) \sim \frac{1}{P_E^2 + M_g^2(P_E^2)}$$

Lattice results \rightarrow an infrared finite propagator (Leinweber et al., hep-lat/0002020 - other possibilities... $1/P_E^4$ ruled out)

One expression for the dynamical gluon mass:

E. Gorbar and A. Natale, Phys.Rev. D61(2000)054012

$$M_g^2(P_E^2) = m_g^2 \theta(m_g^2 - P_E^2) + \frac{m_g^4}{P_E^2} \theta(P_E^2 - m_g^2)$$

with

$$m_g^2 \equiv \left(\frac{34N\pi^2}{9(N^2-1)} \langle \frac{\alpha_s}{\pi} G^{\mu\nu} G_{\mu\nu} \rangle \right)^{1/2}$$

consistent with OPE (Lavelle, PRD44(91)R26)

$$m_g \sim 0.64 \pm 0.20 \text{ GeV}$$

Gluon mass corrections in heavy

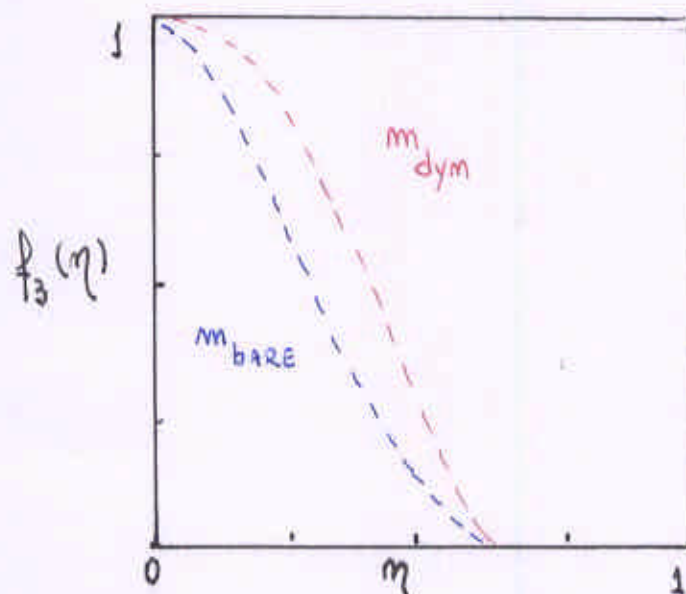
quarkonia decays

A. Mihara and A. Natale, Phys.Lett.B482 (2000) 378

$$R_V = \frac{\Gamma(V \rightarrow ggg)}{\Gamma(V \rightarrow ee)} \rightarrow + \text{gluon mass} \rightarrow f_3(\eta) R_V$$

$$\text{where } \eta \equiv \frac{2m_g}{M_V}$$

$$f_3(\eta) = \frac{\Gamma(V \rightarrow ggg)|_{m_g}}{\Gamma(V \rightarrow ggg)|_{m_g=0}}$$



$$f_3(\eta_{M_{\Upsilon}}) \sim 0.94 \quad , \quad f_3(\eta_{M_{J/\psi}}) \sim 0.47$$

Results

$\Gamma(\Upsilon \rightarrow \text{hadrons})$ is **not modified** by a dynamical gluon mass. Change in $\Gamma(J/\Psi \rightarrow \text{hadrons})$ is consistent with the prediction of Consoli and Field.

Photon spectrum in $J/\Psi \rightarrow \gamma + X$ at $z_\gamma \rightarrow 1$ is totally explained by the gluon mass.

Photon spectrum in $\Upsilon \rightarrow \gamma + X$? (dynamical gluon mass plus soft gluon emission ?)

Dynamical gluon mass also affects the running of α_s at small momentum

$$\alpha_s(Q^2) \rightarrow \frac{12\pi}{(11N - 2n_f) \ln[(Q^2 + \xi m_g^2)/\Lambda^2]}$$

where $\xi \sim 4$. $(\alpha_s^{m_g=0} / \alpha_s^{m_g})$ is a factor 1.6 at the scale m_c . $\alpha_s(m_b) \leftrightarrow \alpha_s(m_c)$?