

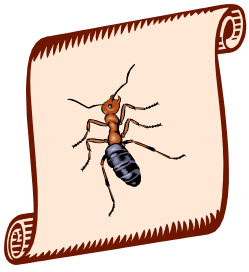
Mini-Review on Extra Dimensions

Greg Landsberg



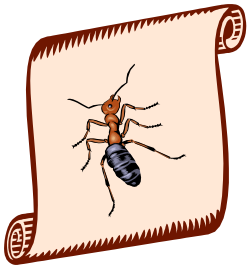
ICHEP 2000

July 28, 2000



Outline

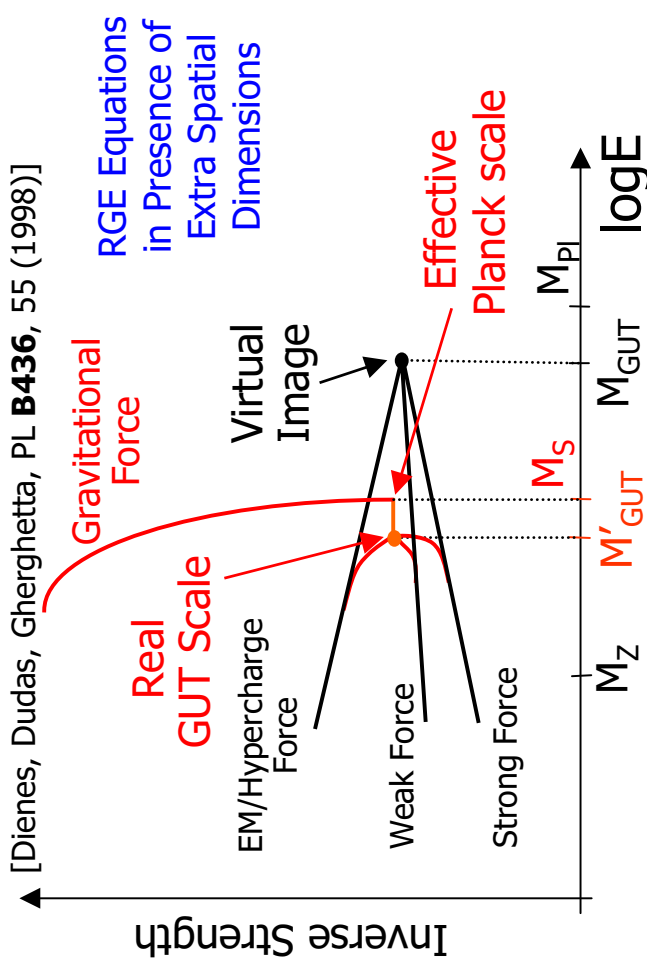
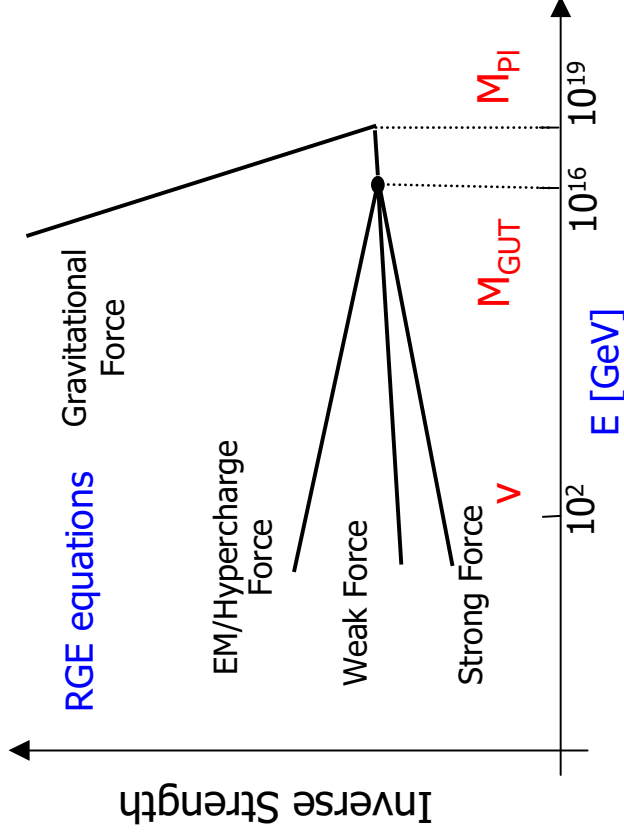
- # Theory of Large Extra Dimensions
- # Current Limits on Large Extra Dimensions
 - # Cosmological Constraints
 - # Gravity at Short Distances
 - # LEP2 Searches for Large Extra Dimensions
 - # Direct Graviton Emission
 - # Virtual Graviton Effects
 - # HERA Searches for Large Extra Dimensions
 - # Tevatron Searches for Large Extra Dimensions
 - # $D\bar{D}$ Search for virtual graviton effects
 - # Looking for direct graviton emission
- # Unusual Signatures for Extra Dimensions
- # Conclusions

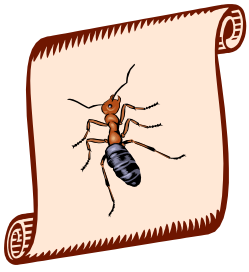


Life Beyond the Standard Model

- ✚ The SM, despite its tremendous success, accommodates, but fails to explain:
 - ✚ EWSB
 - ✚ Origin of fermion masses
 - ✚ CP-violation and baryon asymmetry
- ✚ It also suffers from quadratic divergencies which call for new physics at a scale of the order of Higgs mass, i.e. ~ 1 TeV

- ✚ But: what if there is no other scale, and the SM model is correct up to the Planck scale?
- ✚ Arkani-Hamed, Dimopoulos, Dvali (ADD) (1998): what if the Planck scale is ~ 1 TeV?!
- ✚ What about GUT?



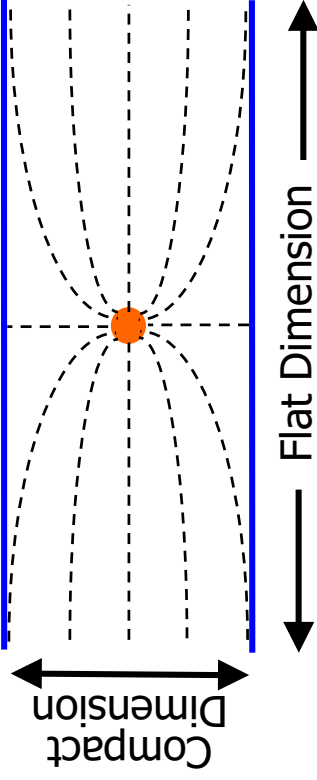


A Crazy Idea? – But it Works!

- But: how to make **gravity strong**?

$$V(r) = \frac{1}{M_{Pl}^2} \frac{m_1 m_2}{r} \rightarrow \frac{1}{(M_{Pl}^{[3+n]})^{n+2}} \frac{m_1 m_2}{r^{n+1}}$$

- Ruled out for **flat extra dimensions**, but has not been ruled out for sufficiently small compactified extra dimensions:



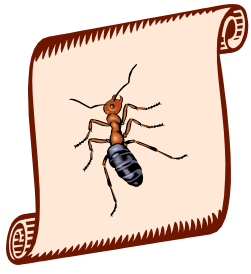
$$V(r) \propto \frac{1}{(M_{Pl}^{[3+n]})^{n+2}} \frac{m_1 m_2}{R^n r} \text{ for } r \gg R$$

$\leftarrow M_S$ – effective Planck Scale

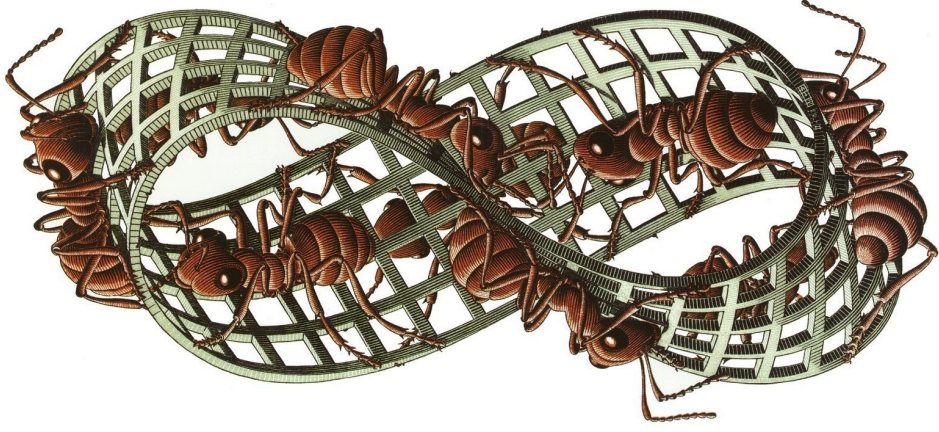
- More precisely, from Gauss's law:

$$R = \frac{1}{2\sqrt{\pi} M_S} \left(\frac{M_{Pl}}{M_S} \right)^{2/n} \propto \begin{cases} 8 \times 10^{12} m, & n=1 \\ 0.7 mm, & n=2 \\ 3 nm, & n=3 \\ 6 \times 10^{-12} m, & n=4 \end{cases}$$

- Amazing as it is, but **no one has tested Newton's law to distances less than ~ 1mm**
- Therefore, **large spatial extra dimensions compactified at a sub-millimeter scale are, in principle, allowed!**



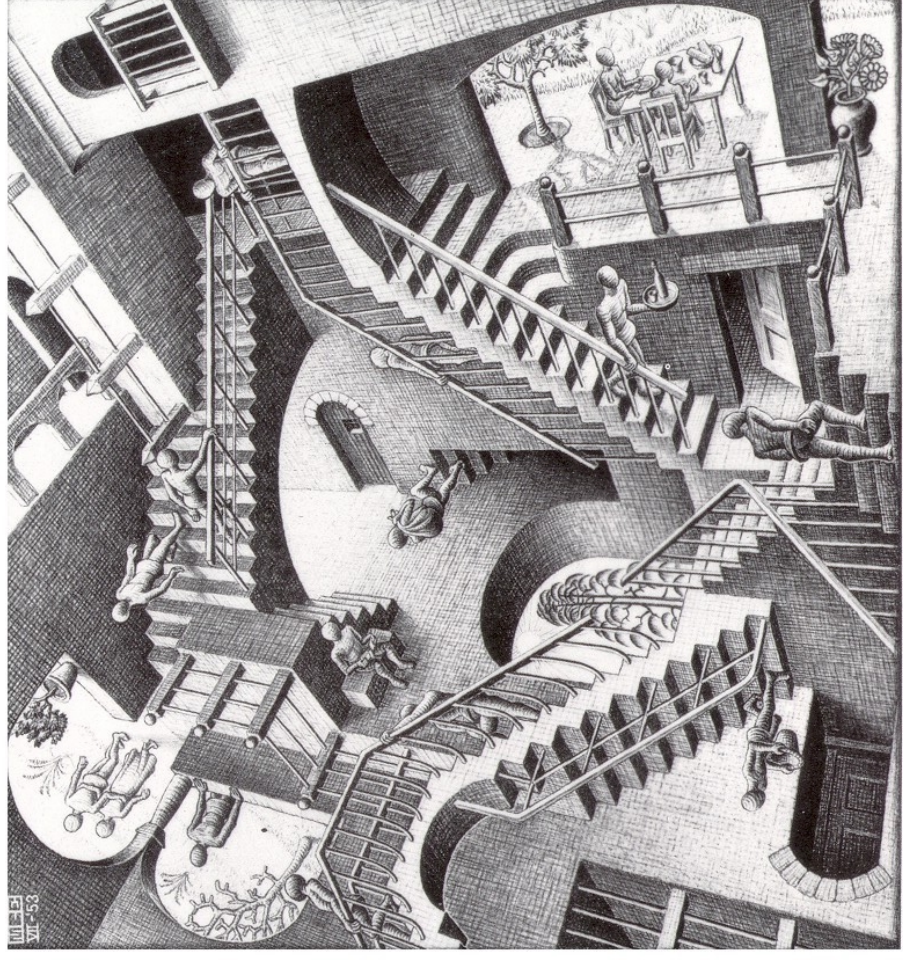
Examples of Compactified Spatial Dimensions



M.C. Escher, Möbius Strip II (1963)

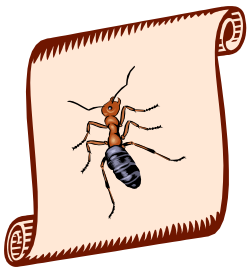
[All M.C. Escher works and texts copyright © Gordon Art B.V., P.O. Box 101, 3740 AC The Netherlands. Used by permission.]

ICHEP 2000, Osaka



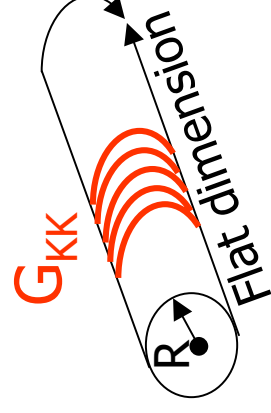
M.C. Escher, Relativity (1953)

Greg Landsberg, Mini-Review on Extra Dimensions



An Importance of Being Compact

- ➦ Compactified dimensions offer a way to **increase tremendously gravitational interaction** due to a large number of the available “winding” modes
- ➦ This tower of excitations is known as **Kaluza-Klein modes**, and such gravitons propagating in the compactified extra dimensions are called Kaluza-Klein gravitons, G_{KK}
- ➦ From the point of view of a 3+1-dimensional space time, the **Kaluza-Klein graviton modes are massive**, with the mass per excitation more $\sim 1/R$
- ➦ Since the mass per excitation mode is so small (e.g. 400 eV for $n = 3$, or 0.2 MeV for $n = 4$), **a very large number of modes can be excited** at high energies

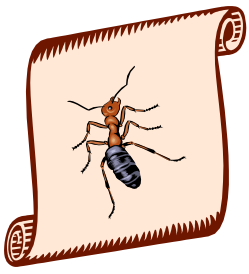


Compactified dimension

$$\phi(x) = \phi(x + 2\pi kR), \quad k = 0, 1, 2, \dots$$

$$M(G_{KK}) = \sqrt{P_x^2} = 2\pi k/R$$

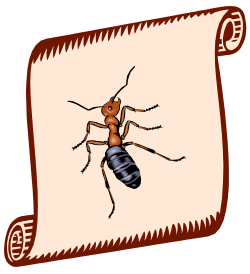
- ➦ Each **Kaluza-Klein graviton mode couples with the gravitational strength**
- ➦ For a large number of modes, **accessible at high energies, gravitational coupling is therefore enhanced drastically**
- ➦ **Low energy** precision measurements **are not sensitive** to the ADD effects



Cosmological Limits on Large Extra Dimensions

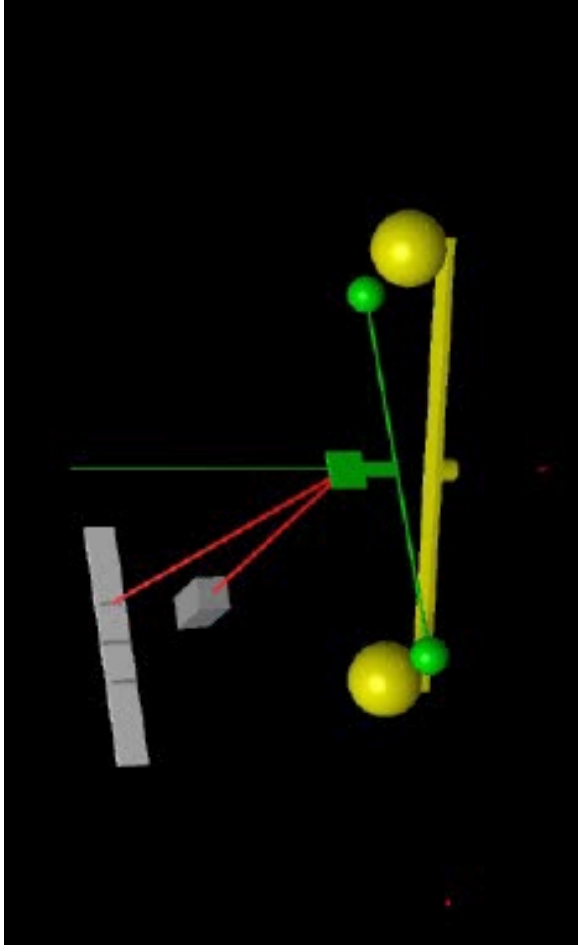
- ✚ Supernova cooling due to the graviton emission
 - ✚ Any new cooling mechanism would decrease the thought-to-be dominant cooling by the neutrino emission
 - ✚ Tightest limits on any additional cooling sources come from the measurement of the SN1987A neutrino flux by the Kamiokande and IMB
 - ✚ Application to the ADD scenario [Cullen, Perelstein, PRL **83**, 268 (1999)]:
 - ✚ $M_S > 30$ TeV ($n=2$)
 - ✚ $M_S > 4$ TeV ($n=3$)

- ✚ Distortion of the cosmic diffuse gamma radiation (CDG) spectrum due to the $G_{KK} \rightarrow \gamma\gamma$ decays
 - ✚ Best CDG measurement come from the COMPTTEL instrument in the 800 KeV - 30 MeV range
 - ✚ Application to the ADD scenario [Hall, Smith, PRD **60**, 085008 (1999)]:
 - ✚ $M_S > 100$ TeV ($n=2$)
 - ✚ $M_S > 5$ TeV ($n=3$)
 - ✚ **Caveat:** there are many known (and unknown!) uncertainties, so the cosmological bounds are reliable only as an order of magnitude estimate
 - ✚ Still, $n=2$ seems to be excluded



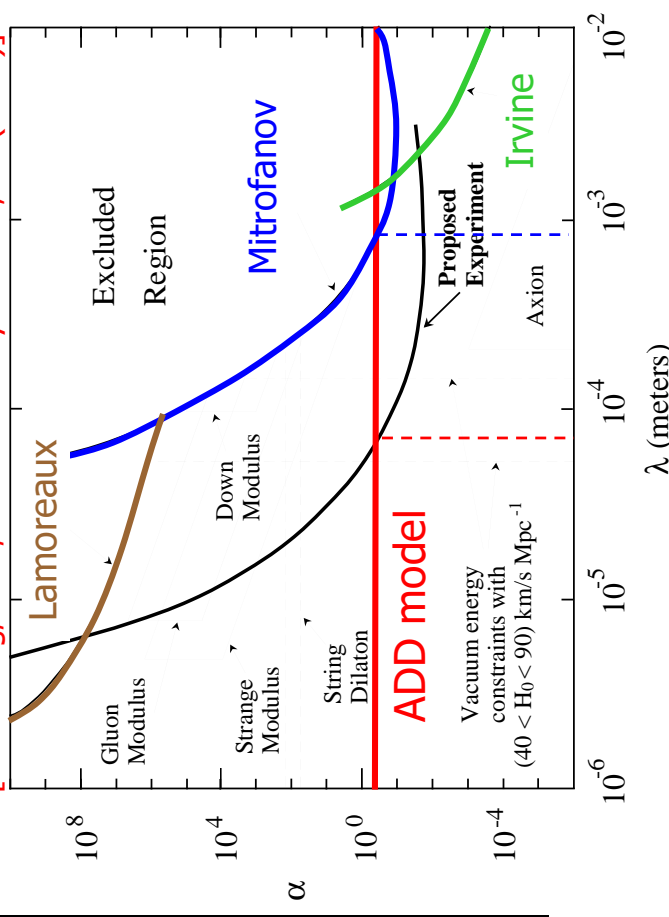
Current Limits from Gravitational Experiments

- 1798: Cavendish experiment (torsion balance)

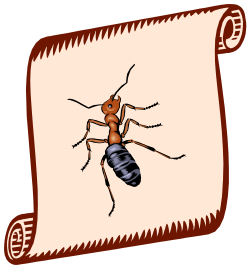


- Mid-1970-ies: a number of Cavendish-type experiments searching for the “fifth forth” via deviations from Newton’s law
- Sensitivity vanishes quickly for distances less than 1 mm
- Major background: Van der Waals and Casimir forces

Status of short-range gravity experiments
[From Long/Chan/Price Nucl. Phys. **B539**, 23 (1999)]



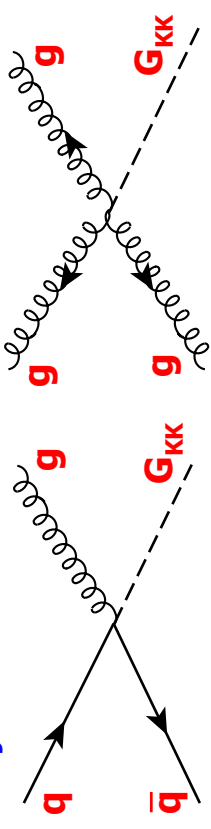
- Best sub-millimeter results are from 1997 Lamoreaux experiment [PRL **78**, 5 (1997)] to measure the Casimir force
- New preliminary result from Adelberger et al. (APS 2000 Meeting) indicates that $n=2$ is ruled out



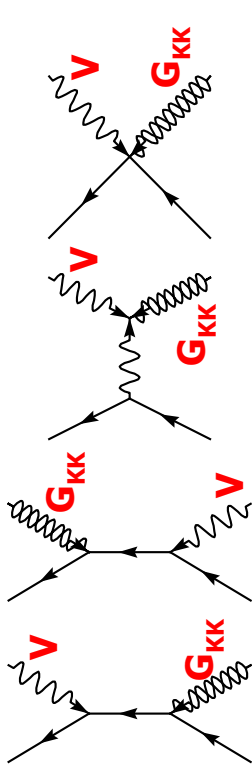
Collider Signatures for Large Extra Dimensions

- Kaluza-Klein gravitons couple to the momentum tensor, and therefore contribute to most of the SM processes
- For Feynman rules for G_{KK} see:
 - Han, Lykken, Zhang, PR **D59**, 105006 (1999)
 - Giudice, Rattazzi, Wells, Nucl. Phys. **B544**, 3 (1999)
- Since graviton can propagate in the bulk, energy and momentum are not conserved in the G_{KK} emission from the point of view of our 3+1 space-time
- Since the spin 2 graviton in generally has a bulk momentum component, its spin from the point of view of our brane can appear as 0, 1, or 2
- Depending on whether the G_{KK} leaves our world or remains virtual, the collider signatures include single photons/Z/jets with missing E_T or fermion/vector boson pair production

Real Graviton Emission
Monojets at hadron colliders

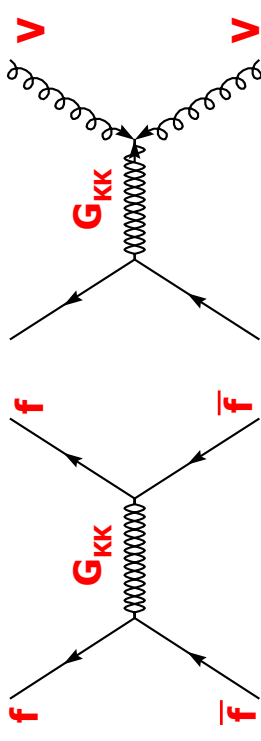


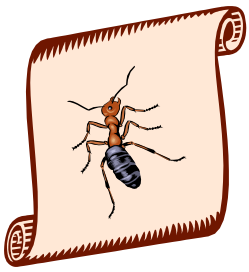
Single VB at hadron or e^+e^- colliders



Virtual Graviton Emission

Fermion or VB pairs at hadron or e^+e^- colliders





LEP2 Searches for Direct Graviton Emission - I

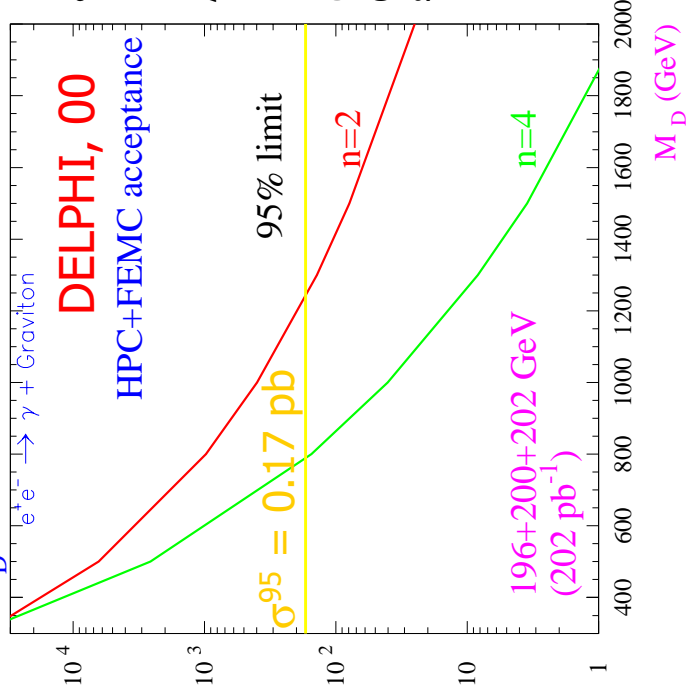
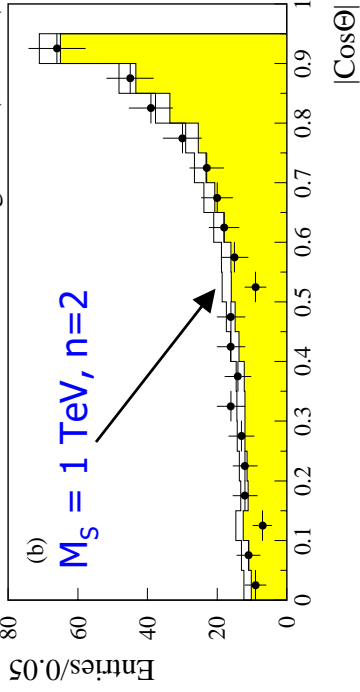
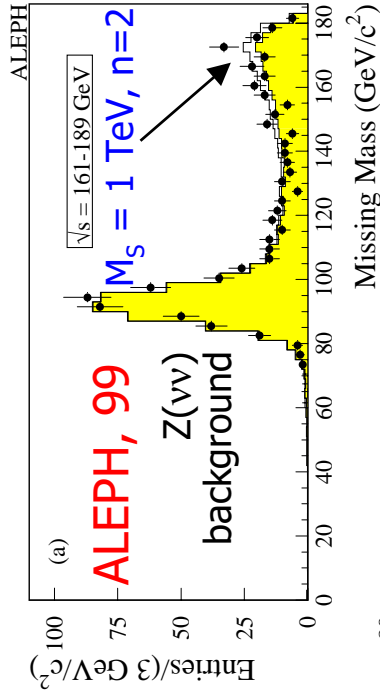
- $e^+e^- \rightarrow \gamma G_{KK}$
- Photon + ME_T signature
- "Recycling" of the GMSB analyses
- ALEPH (2D-fit), DELPHI, L3 (x), OPAL (event counting)

$$\frac{d^2\sigma}{dxdz} = \frac{\alpha}{32s} \frac{\pi^{\frac{n}{2}}}{\Gamma(\frac{n}{2})} \left(\frac{\sqrt{s}}{M_S}\right)^{n+2} f(x, z), \quad x = \frac{2E_\gamma}{\sqrt{s}}, \quad z = \cos\theta$$

$$f(x, z) = \frac{2(1-x)^{\frac{n-1}{2}}}{x(1-z^2)} \left[(2-x)^2 (1-x+x^2) - 3x^2(1-x)z^2 - x^4z^4 \right]$$

$M_D > 1250$ GeV for $n=2$

$M_D > 792$ GeV for $n=4$



Theory:

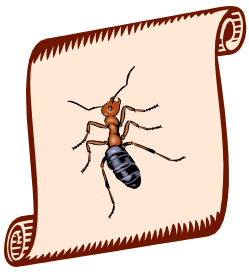
[Giudice, Rattazzi, Wells, Nucl. Phys. **B544**, 3 (1999) and corrected version: hep-ph/9811291]

Experiment:

ALEPH-CONF-2000-005
DELPHI 2000 CONF 344
L3: Phys. Lett. **B470**, 268 (1999)
OPAL: CERN-EP-2000-050, submitted to Eur. Phys. J C

Results:

$M_S > 1.3-0.6$ TeV for $n=2-6$ (DELPHI)
ALEPH, L3, OPAL – slightly worse

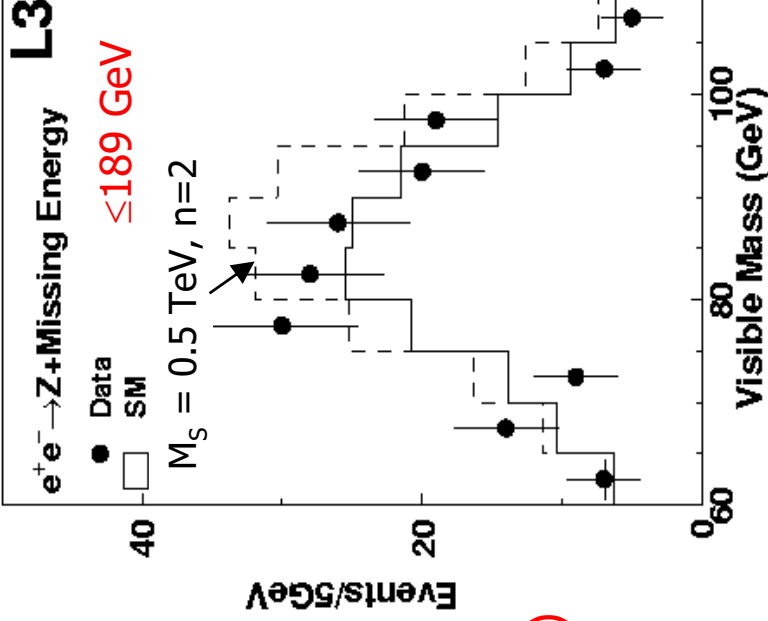


LEP2 Searches for Direct Graviton Emission - II

- $e^+e^- \rightarrow ZG_{KK}$
- $Z(jj) + ME_T$ signature
- "Recycling" of the invisible Higgs analyses
- ALEPH: $Z(jj)G$, 184 GeV, total cross section method
- L3: $Z(jj)G$, 189 GeV, increased sensitivity via analysis of the visible mass distribution
- OPAL: 189 GeV
- $M_S > 0.35\text{-}0.12$ TeV (ALEPH) for $n = 2\text{-}6$
- $M_S > 0.60\text{-}0.21$ TeV (L3) for $n = 2\text{-}6$

$$\frac{\Gamma(Z \rightarrow \bar{f}f G)}{\Gamma(Z \rightarrow \bar{f}f)} = \frac{1}{4\pi} \frac{1}{3(n+2)} \left(\frac{M_Z}{M_S}\right)^{n+2} I$$

$$I = \frac{\pi^{\frac{n-2}{2}}}{\Gamma(\frac{n}{2})} \int_0^1 dx \int_0^{(1-\sqrt{x})^2} dy \frac{y^{\frac{n-2}{2}} (12x+A)\sqrt{A}}{6(1-x)^2}, \quad A = (1-x-y)^2 - 4xy$$

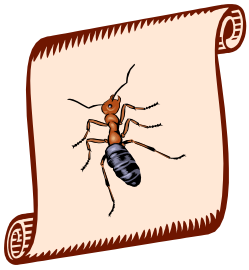


Theory:

[Balazs, Dicus, He, Repko, Yuan, Phys. Rev. Lett. **83**, 2112 (1999) – width ratio]
 [Cheung, Keung, Phys. Rev. **D60**, 112003 (1999) – mass distribution]

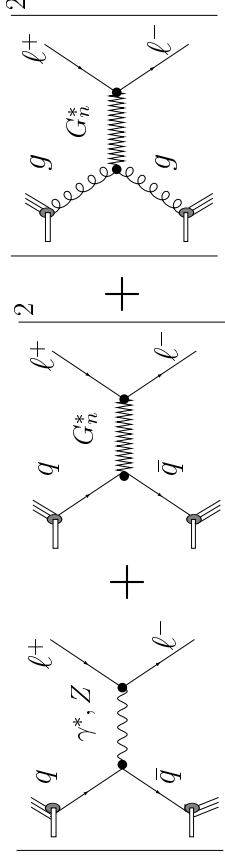
Experiment:

ALEPH-CONF-99-027
 L3: Phys. Lett. **B470**, 281 (1999)



Virtual Graviton Effects

- In the case of pair production via virtual graviton, gravity effects interfere with the SM (e.g., t^+t^- at hadron colliders):



- Therefore, production cross section has three terms: SM, interference, and direct gravity effects
- The sum in KK states is divergent in the effective theory, so in order to calculate the cross sections, an explicit cut-off is required
- An expected value of the cut-off is $\approx M_S$, as this is the scale at which the effective theory breaks down, and the string theory needs to be used to calculate production

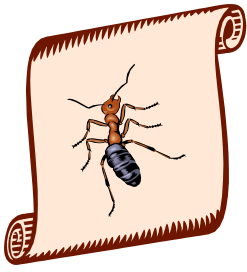
- Unfortunately, a number of similar papers calculating the virtual graviton effects appeared simultaneously

Hence, there are three major conventions on how to write the effective Lagrangian:

- Hewett, Phys. Rev. Lett. **82**, 4765 (1999)
- Giudice, Rattazzi, Wells, Nucl. Phys. **B544**, 3 (1999); revised version, hep-ph/9811291
- Han, Lykken, Zhang, Phys. Rev. **D59**, 105006 (1999); revised version, hep-ph/9811350

- Fortunately (after a lot of discussions and revisions) all three conventions turned out to be completely equivalent and only the definitions of M_S are different:

$$\frac{d^2 \sigma}{d \cos \theta^* dM} = \frac{d^2 \sigma_{SM}}{d \cos \theta^* dM} + \frac{a(n)}{M_S^4} f_1(\cos \theta^*, M) + \frac{b(n)}{M_S^8} f_2(\cos \theta^*, M)$$



Hewett, GRW, and HLZ Formalisms

- Hewett:** neither sign of the interference nor the dependence on the number of extra dimensions is known; therefore the **interference term is $\sim \lambda/M_S^4$ (Hewett)**, where λ is of order 1; numerically uses $\lambda = \pm 1$
- GRW:** sign of the interference is fixed, but the dependence on the number of extra dimensions is unknown; therefore the **interference term is $\sim 1/\Lambda_T^4$** (where Λ_T is their notation for M_S)
- HLZ:** not only the sign of interference is fixed, but the n-dependence can be calculated in the effective theory; thus the **interference term is $\sim F/M_S^4$ (HLZ)**, where **F** reflects the dependence on the number of extra dimensions:

$$F = \begin{cases} \log\left(\frac{M_S^2}{s}\right), & n = 2 \\ \frac{2}{n-2}, & n > 2 \end{cases}$$

- Correspondence** between the three formalisms:

$$M_S(\text{Hewett})_{\lambda=\pm 1} \equiv \sqrt[4]{\frac{2}{\pi} \Lambda_T}(\text{GRW})$$

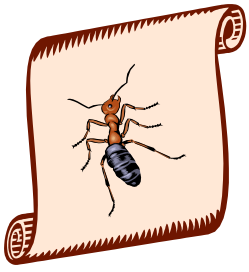
$$\frac{\lambda}{M_S^4(\text{Hewett})} = \frac{\pi}{2} \frac{F}{M_S^4(\text{HLZ})}$$

$$\frac{1}{\Lambda_T^4(\text{GRW})} = \frac{F}{M_S^4(\text{HLZ})}$$

- Rule of thumb:**

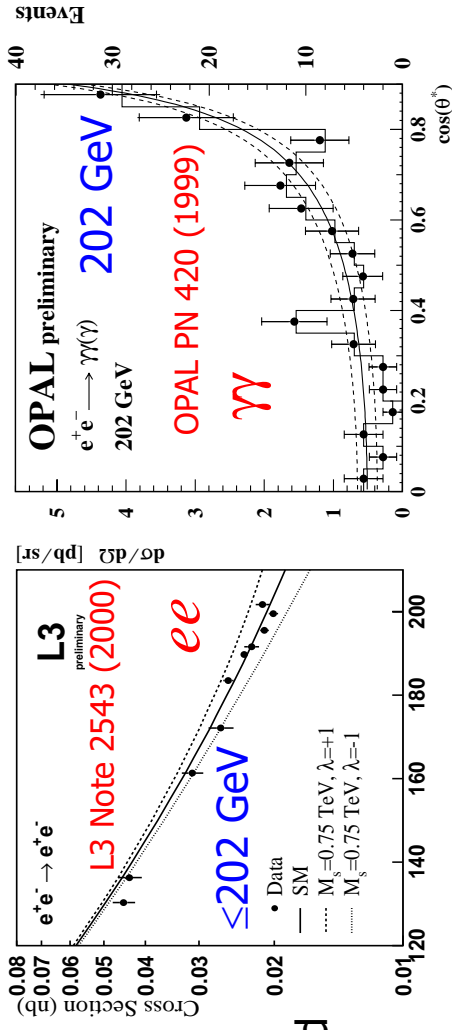
$$M_S(\text{Hewett})_{\lambda=\pm 1} \approx M_S(\text{HLZ})_{n=5}$$

$$\Lambda_T(\text{GRW}) = M_S(\text{HLZ})_{n=4}$$

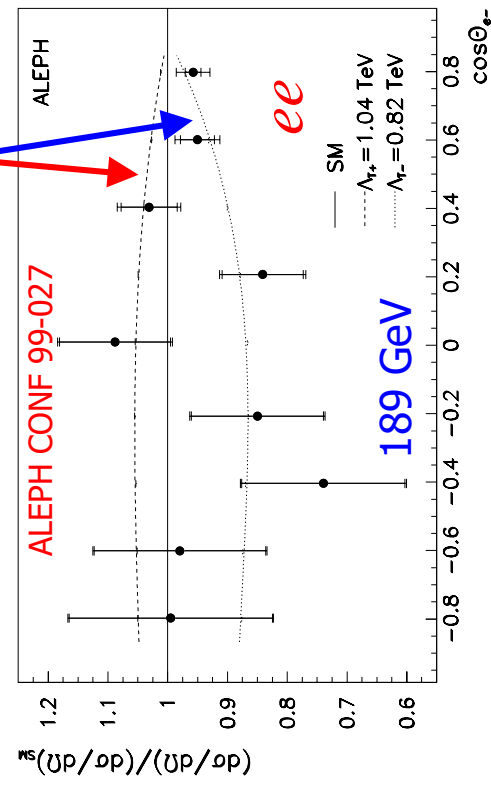


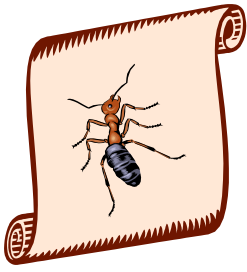
LEP2 Searches for Virtual Graviton Effects

- LEP2 Collaborations looked at **difermion** and **diboson** production due to the G_{KK} exchange
- Unfortunately, **different formalisms were used by different collaborations**, and sometimes even within a collaboration, which makes results hard to compare and combine
- Internal inconsistency** could affect some of the **combined limits**
- Most sensitive channels** are:
 - Dielectron s-channel production and Bhabha scattering
 - Diphoton production
- Limits on M_S (Hewett) \sim **0.8-1.2 TeV**
- Bibliography:
 - ALEPH: CONF 99-027, 2000-005, 2000-030
 - DELPHI: CONF 355, 363 (2000)
 - L3: PL **B464**, 135; **B470**, 281 (1999); Notes 2543, 2590 (2000)
 - OPAL: CERN-EP/99-097, PN 420 (1999)



N.B. All LEP Collaborations considered both interference signs





LEP2 Searches for Virtual Graviton Effects - WV

- $e^+e^- \rightarrow WW/ZZ$
- Recycle WW cross section and anomalous $ZZ\gamma$ couplings analyses
- L3 used angular distributions (WW) and mass variables (ZZ) to set limits

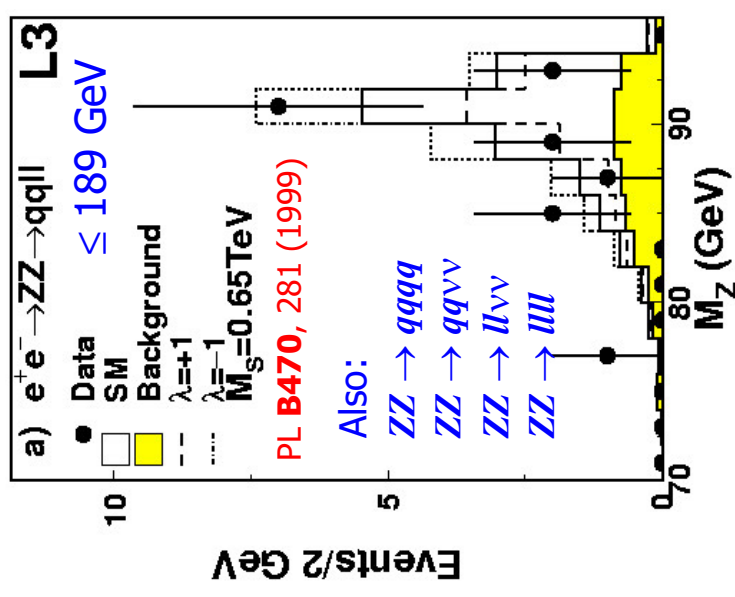
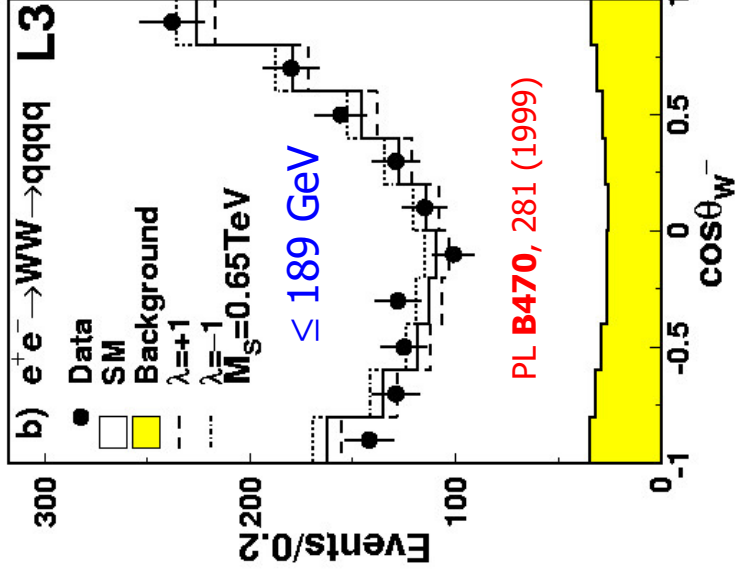
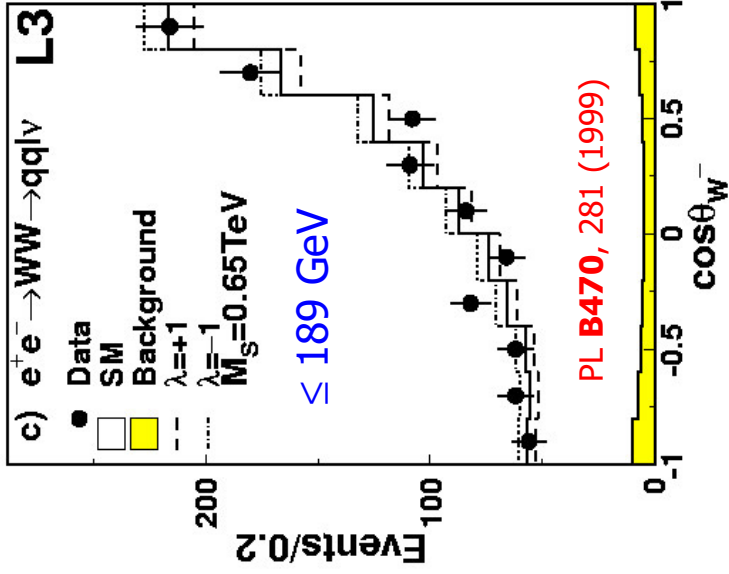
Theory:

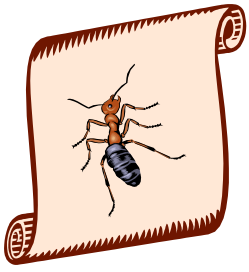
[Agashe, Deshpande, Phys. Lett. **B456**, 60 (1999)]

$$M_S^{AD} \Big|_{\lambda=-1} \equiv M_S^{Hewett} \Big|_{\lambda=+1}$$

AD convention is equivalent to Hewett's with a flipped sign of λ

$M_S > 520\text{-}650$ GeV (WW); $M_S > 460\text{-}470$ GeV (ZZ)



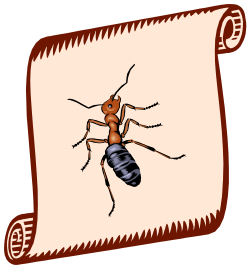


LEP2 Lower 95% CL M_S (Hewlett) Limits (TeV)

Experiment	$e^+e^- \rightarrow \gamma G$						$e^+e^- \rightarrow ZG$						Color coding		
	n=2	n=3	n=4	n=5	n=6	n=6	n=2	n=3	n=4	n=5	n=6	n=6	n=6	$\lambda=-1$	$\lambda=+1$
ALEPH	1.10	0.86	0.70	0.60	0.52	0.52	0.35	0.22	0.17	0.14	0.12			≤ 184 GeV	
DELPHI	1.25	0.97	0.79	0.68	0.59									≤ 189 GeV	
L3	1.02	0.81	0.67	0.58	0.51	0.51	0.60	0.38	0.29	0.24	0.21			≤ 202 GeV	
OPAL	1.09	0.86	0.71	0.61	0.53	0.53								$\lambda=-1$	$\lambda=+1$

Virtual Graviton Exchange

Experiment	e^+e^-	$\mu^+\mu^-$	$\tau^+\tau^-$	qq	ff	$\gamma\gamma$	WW	ZZ	Combined
ALEPH (Λ_T)	1.17 0.91	0.73 0.75	0.67 0.69	0.59/0.64 0.55/0.55 (bb)	1.17 0.94	0.91 0.92			0.84/1.12 (<189) $M_S > 0.75/1.00$
DELPHI		0.59 0.73	0.56 0.65		0.60 0.76	0.69 0.71			0.60/0.76 (ff) (<202)
L3	0.91 0.99	0.56 0.69	0.58 0.54	0.49 0.49	0.84 1.00	0.80 0.79	0.68 0.79	1.2 1.2	1.3/1.2 (<202)
OPAL		0.63 0.60	0.50 0.63		0.61 0.68	0.63 0.64			0.61/0.68 (ff) (<189)



HERA Search for Virtual Graviton Effects

$$e^\pm p \rightarrow e^\pm p$$

- ✚ t-channel exchange, similar to Bhabha scattering diagrams; based on the GRW formalism (set limits on Λ_T , but call it M_S)
- ✚ Usual SM, Z/γ^* interference, and direct G_{KK} terms
- ✚ Analysis method: fit to the $d\sigma/dQ^2$ distribution (H1)
- ✚ Current combined limits: $\Lambda_T > 0.67/0.73$ TeV ($M_S > 0.60/0.65$ TeV)
- ✚ Expected sensitivity up to 1 TeV with the ultimate HERA data set

$$\frac{d\sigma(e^\pm q \rightarrow e^\pm q)}{dt} = \frac{d\sigma_{SM}}{dt} + \frac{d\sigma_G}{dt} + \frac{d\sigma_{\gamma G}}{dt} + \frac{d\sigma_{ZG}}{dt},$$

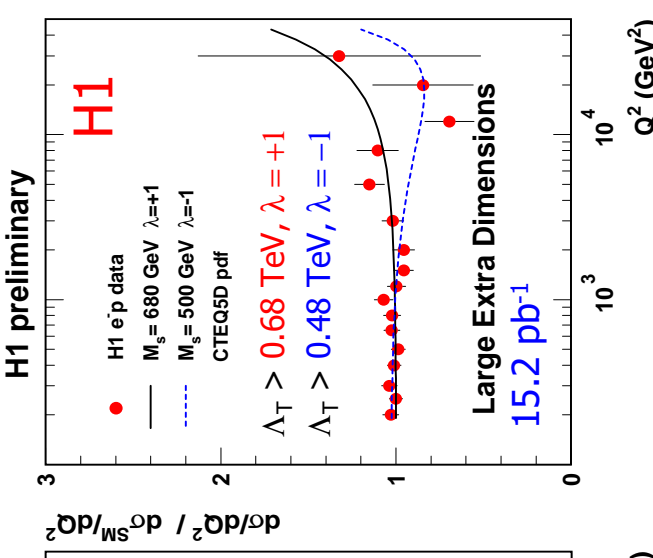
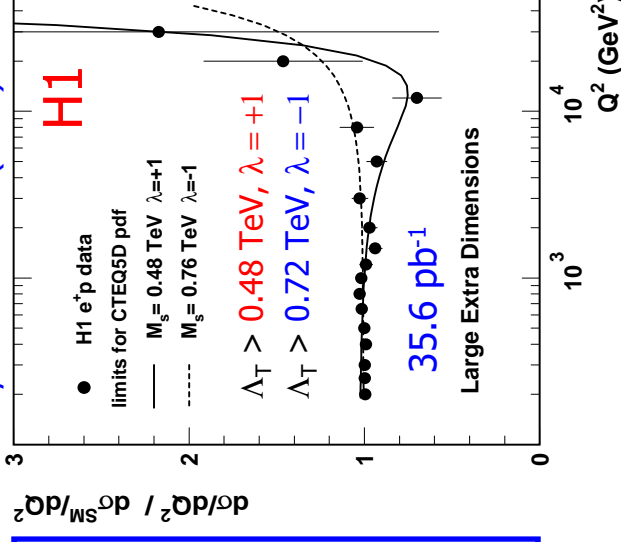
$$\frac{d\sigma_G}{dt} = \frac{\pi \lambda^2}{32 M_S^8 s^2} (32u^4 + 64u^3 t + 42u^2 t^2 + 10u t^3 + t^4),$$

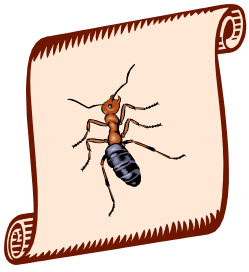
$$\frac{d\sigma_{\gamma G}}{dt} = \frac{\pi \lambda}{2 M_S^4 s^2} \frac{\alpha e_q (2u+t)}{t},$$

$$\frac{d\sigma_{ZG}}{dt} = \frac{\pi \lambda}{2 M_S^4 s^2} \sin^2 2\theta_W \left(v_e v_q \frac{(2u+t)^3}{t - m_Z^2} - a_e a_q \frac{t(6u^2 + 6ut + t^2)}{t - m_Z^2} \right),$$

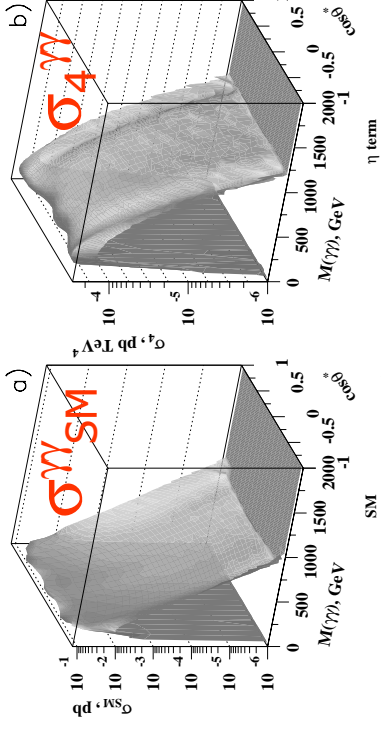
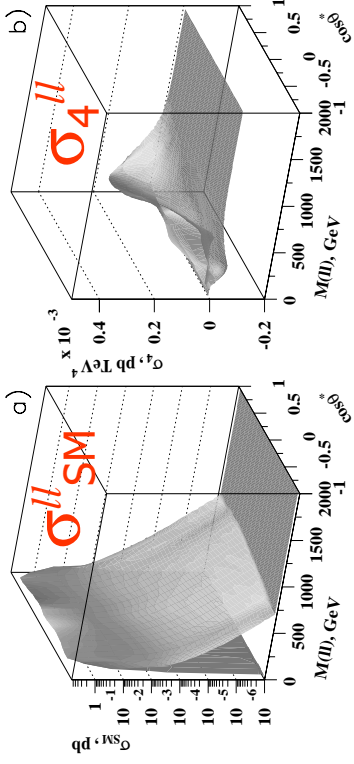
$$\frac{d\sigma(e^\pm g \rightarrow e^\pm g)}{dt} = \frac{\pi \lambda^2}{8 M_S^8 s^2} u (2u^3 + 4u^2 t + 3ut^2 + t^3)$$

Phys. Lett. **B479** (2000) 358





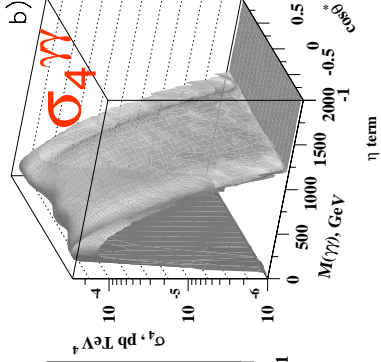
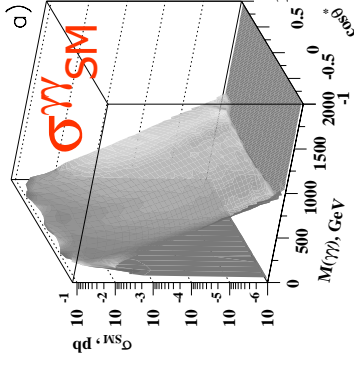
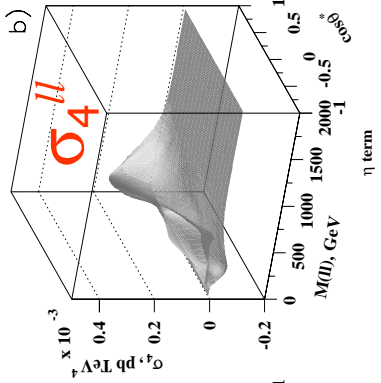
Tevatron Searches for Virtual Graviton Effects



$$\eta = \frac{F}{M_S^4}$$

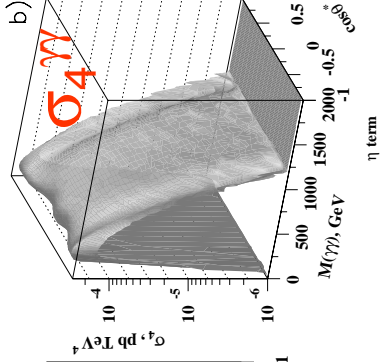
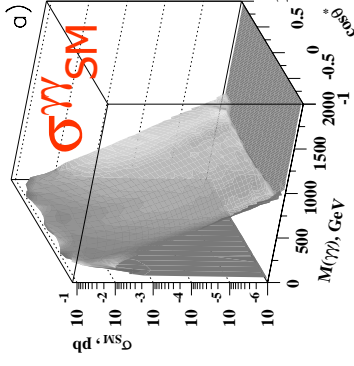
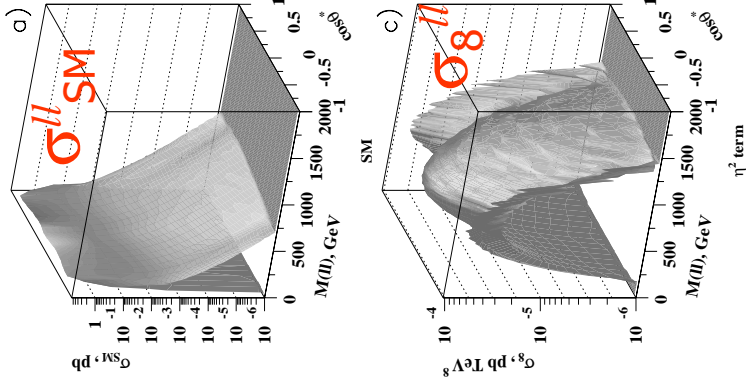
$$F = \begin{cases} \log\left(\frac{M_S^2}{2}\right), & n = 2 \\ \frac{1}{n-2}, & n > 2 \end{cases}$$

$$[\eta] = \text{TeV}^{-4}$$



$$\sigma_{\eta}^{ll} = \sigma_{SM}^{ll} + \sigma_4^{ll} \eta + \sigma_8^{ll} \eta^2$$

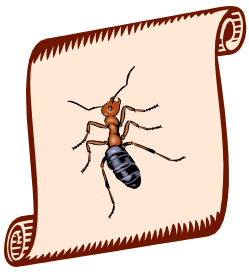
$$\sigma_{\eta}^{\mathcal{W}} = \sigma_{SM}^{\mathcal{W}} + \sigma_4^{\mathcal{W}} \eta + \sigma_8^{\mathcal{W}} \eta^2$$



[Cheung, GL – hep-ph/9909218, to appear in PRD]

- Use **invariant mass** and **scattering angle** in the c.o.m. frame to **maximize sensitivity**
- Parameterize cross section as a **bilinear form in scale η** (works for any $n > 2$)
- Note the **asymmetry** of the interference term, σ_{4l} , for ll production

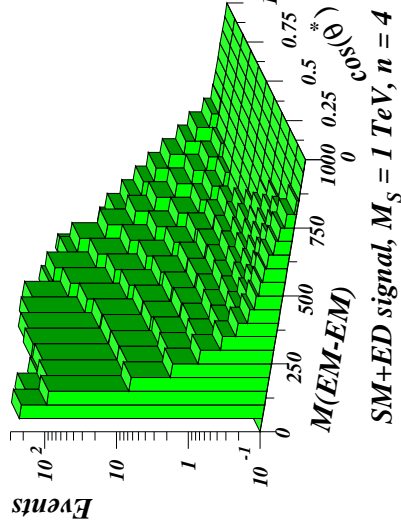
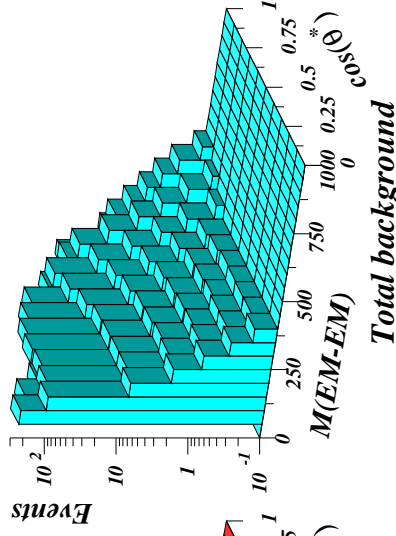
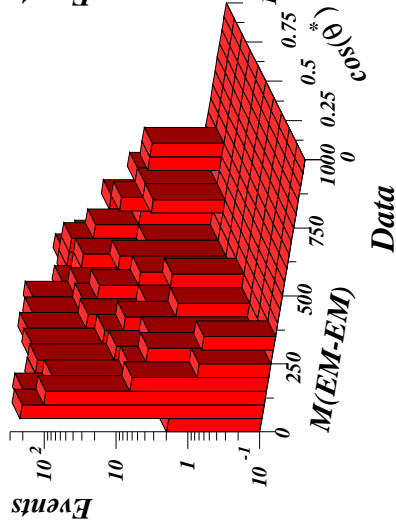
- Use **Bayesian fit** to the data (real one or MC one) to get the best estimate of η
- Diphoton** channel is considerably **more sensitive** than the **dilepton** one



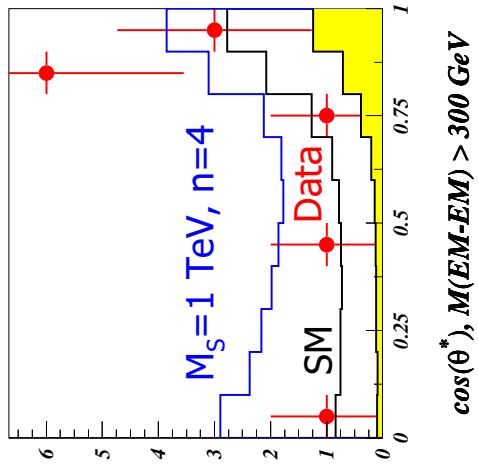
DØ Search Using Dielectrons and Diphotons

Comparison of the data and the SM predictions

DØ Preliminary, Run I, 127 pb⁻¹



SM+ED signal, $M_S = 1 \text{ TeV}, n = 4$

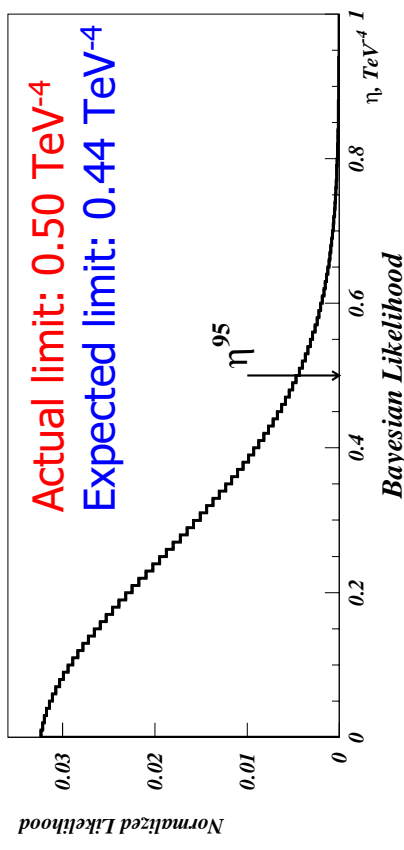


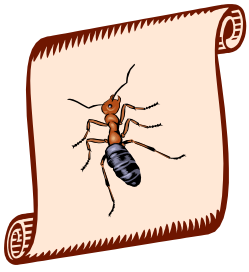
$\cos(\theta^*), M(\text{EM-EM}) > 300 \text{ GeV}$

- Combine dielectron and diphoton events to avoid inefficiency due to tracking and conversions

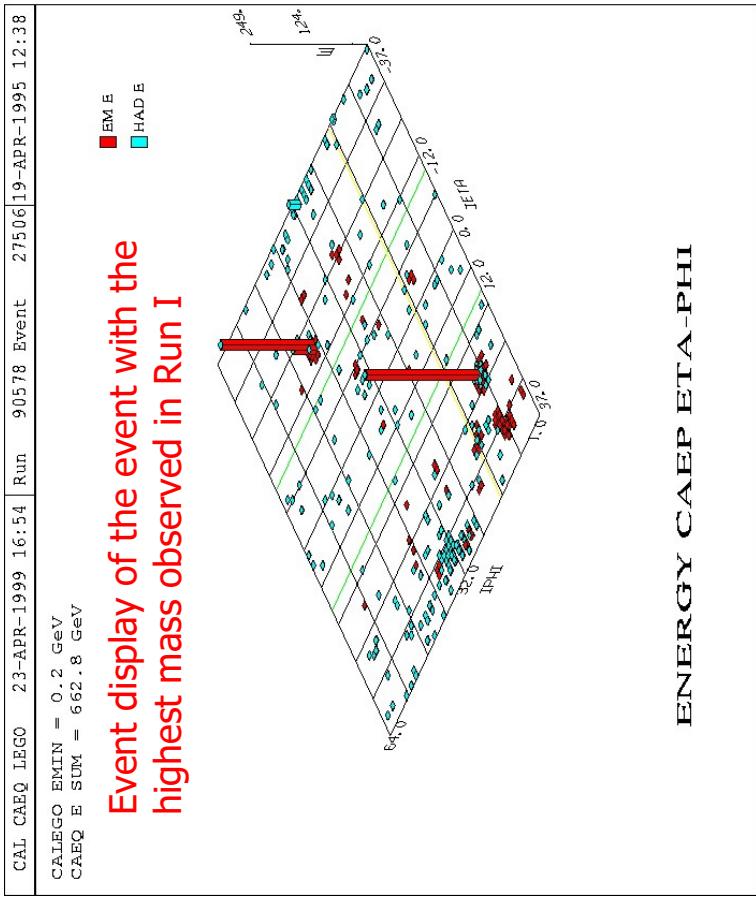
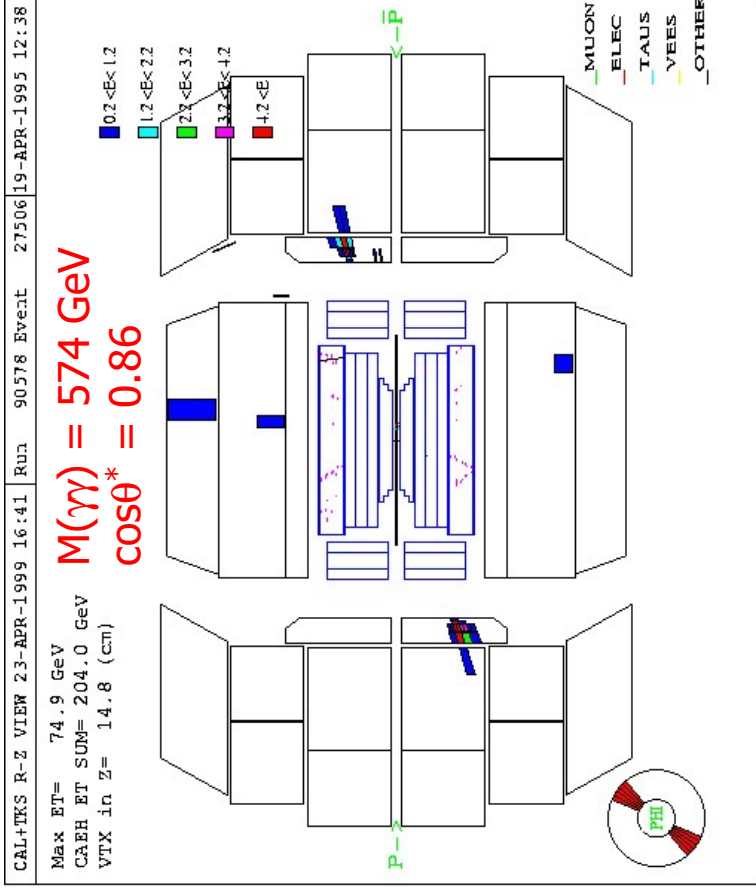
- Since the dominant background is irreducible DY and diphoton production, release ID cuts to maximize sensitivity

- No excess of events is seen at high masses and low scattering angles, where the signal is expected to exhibit itself



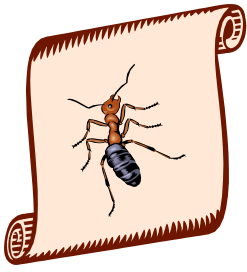


High-Mass Candidate Events



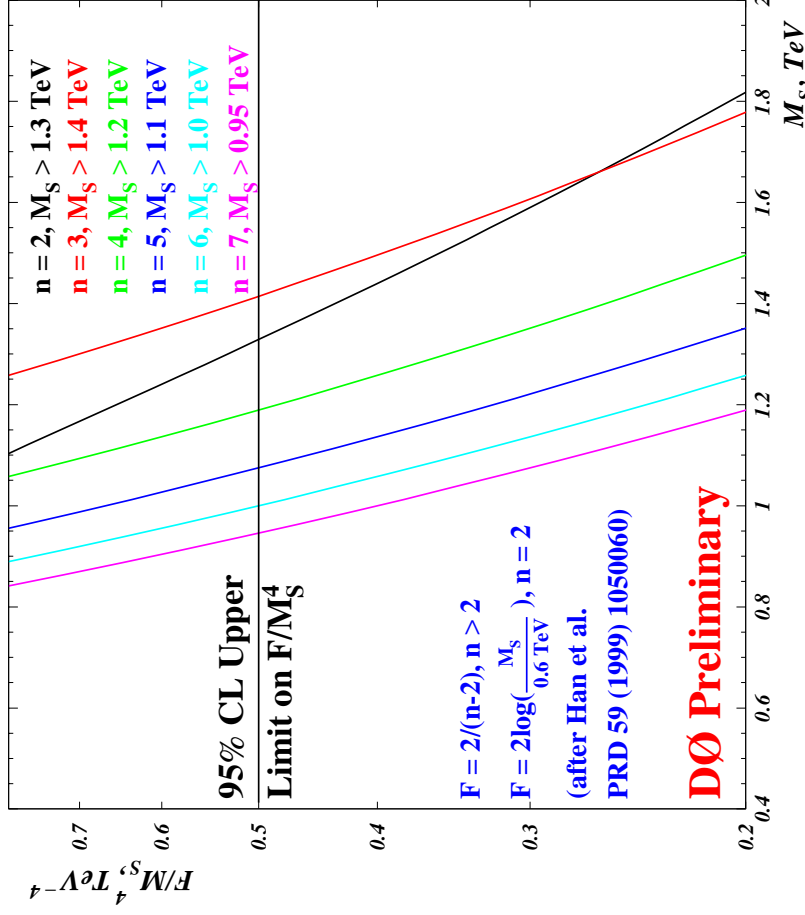
Parameters of the two high-mass candidate events:

Run	Event	Z_{vtx}	M_{E_T}	Type	E_T^1	E_T^2	η_1	η_2	M	$\cos\theta^*$	N_{jet}	$P_T\text{-kick}$
90578	27506	3.6 cm	15 GeV	γ	81 GeV	81 GeV	1.98	-1.91	575 GeV	0.86	0	11.7 GeV
84582	11674	-34 cm	15 GeV	ee	134 GeV	132 GeV	0.99	-1.59	520 GeV	0.84	0	18.8 GeV



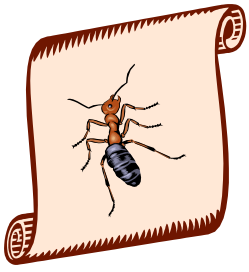
DØ Limits on Large Extra Dimensions

Limits on Large Spatial Extra Dimensions



- Sensitivity is limited by statistics; it will double (in terms of M_S) in Run IIA (2 fb^{-1}), and triple in Run IIB (20 fb^{-1})

- For $n > 2$ M_S limits can be obtained directly from η limits
- For $n = 2$, use averages for gravity contribution ($\langle \hat{s} \rangle = 0.36 \text{ TeV}^2$, see hep-ph/9909218)
- Translate limits in the Hewett and GRW frameworks for easy comparison with other experiments:
 - $M_S(\text{Hewett}) > 1.1 \text{ TeV}$ and 1.0 TeV ($\lambda = -1$)
 - $\Lambda_{\text{T}}(\text{GRW}) > 1.2 \text{ TeV}$
- This limits are comparable with the latest limits from LEP2
- They are complementary to those from LEP2, as they probe different range of energies
- Looking forward for limits from the CDF DY analysis ($M_S \sim 0.9\text{-}1.0 \text{ TeV}$), utilizing the same technique
- Significant improvement is expected by combining the results of the two Tevatron experiments



Tevatron: Real Graviton Emission

- $q\bar{q} \rightarrow gG_{KK}$ (dominant channel)
- **jets + ME_T** final state
- **$Z(\nu\nu)$ +jets** is irreducible **background**
- Important **instrumental backgrounds** from jet mismeasurement, cosmics, etc.
- Both **CDF** and **DØ** are pursuing this search

Theory:

[Giudice, Rattazzi, Wells, Nucl. Phys. **B544**, 3 (1999) and corrected version, hep-ph/9811291]

[Mirabelli, Perelstein, Peskin, PRL **82**, 2236 (1999)]

$$\frac{d\sigma}{dt}(q\bar{q} \rightarrow gG) = \frac{\alpha_s}{36} \frac{1}{sM_s^2} F_1\left(\frac{t}{s}, \frac{m^2}{s}\right)$$

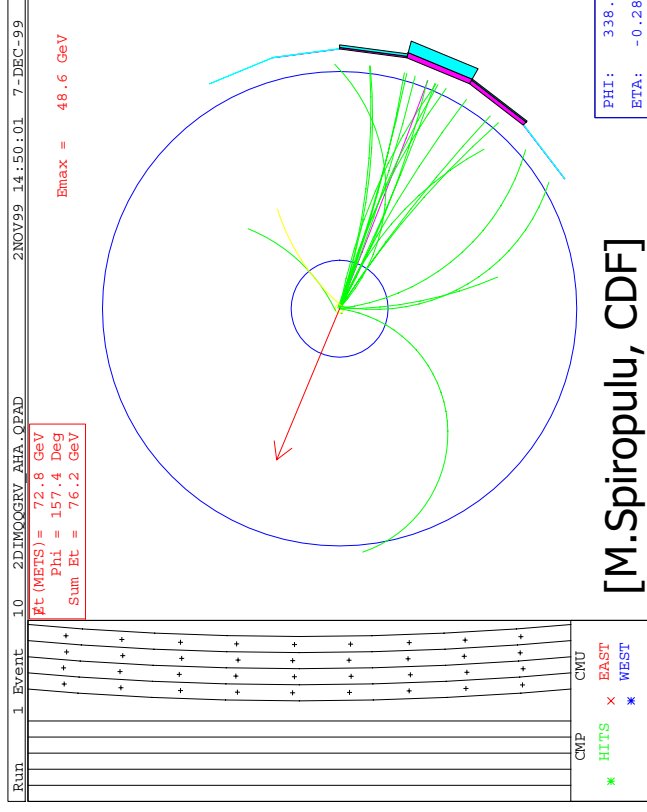
$$F_1(x, y) = \frac{1}{x(y-1-x)} \left[-4x(1+x)(1+2x+2x^2) + y(1+6x+18x^2+16x^3) - 6y^2x(1+2x) + y^3(1+4x) \right]$$

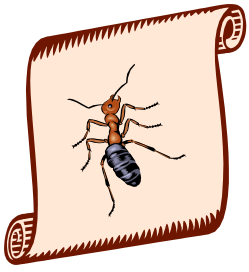
Tevatron Run I/II reach, CDF+DØ [Giudice et al.]

n	M_s reach, Run I	M_s reach, Run II
2	1100 GeV	1400 GeV
3	950 GeV	1150 GeV
4	850 GeV	1000 GeV
5	700 GeV	900 GeV

Note that non-perturbative effects could be important at high n

Note that this sensitivity estimate is probably optimistic, as it does not take into account copious instrumental backgrounds

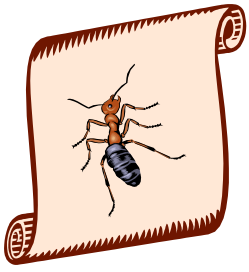




Black Hole Production

- Once the **c.m. energy exceeds the compactification scale**, M_S , a critical energy density is achieved and the black hole is formed
- Not to worry** about the Earth being sucked into such a black hole; they should be constantly formed by cosmic rays
- The temperature of such a black hole is:
 $T = M_{\text{Pl}}^2/M \rightarrow M_S^2/M \times O(M/M_S) \sim M_S$
- For $M_S \sim T = 1 \text{ TeV}$, the **black body spectrum peaks at 250 GeV**, and therefore the BH technically evaporates by emitting a single energetic photon – not quite a black body!
- Moreover, the **lifetime** of such a black hole is only $\sim 10^{-29} \text{ s}$
- The **Swartzchild radius** of such a black hole is $\sim 1/M_S$, i.e. it's \sim **de Broglie wavelength**; it's not clear of one could even consider such an object as a bound state

- Other possibility is **evaporation in the bulk via G_{KK}** , in which case the signature is a **deficit of high-s events**
 - At a hadron collider it's **easy to tweak p.d.f.** to account for such a deficit
 - At a lepton collider it's **hard to establish that the beams have not missed each other** in one of the better established spatial dimensions
- Interesting possibility for a black hole is to have a **color 'hair' that holds it to our brane**; if the color quantum number is conserved, the black hole could be metastable and live seconds or even days before it decays in a large number of hadrons
 - Look for **events not in time with the accelerator clock with such a distinct signature** (Dvali, GL, Matchev)






Gauge Boson Excitations

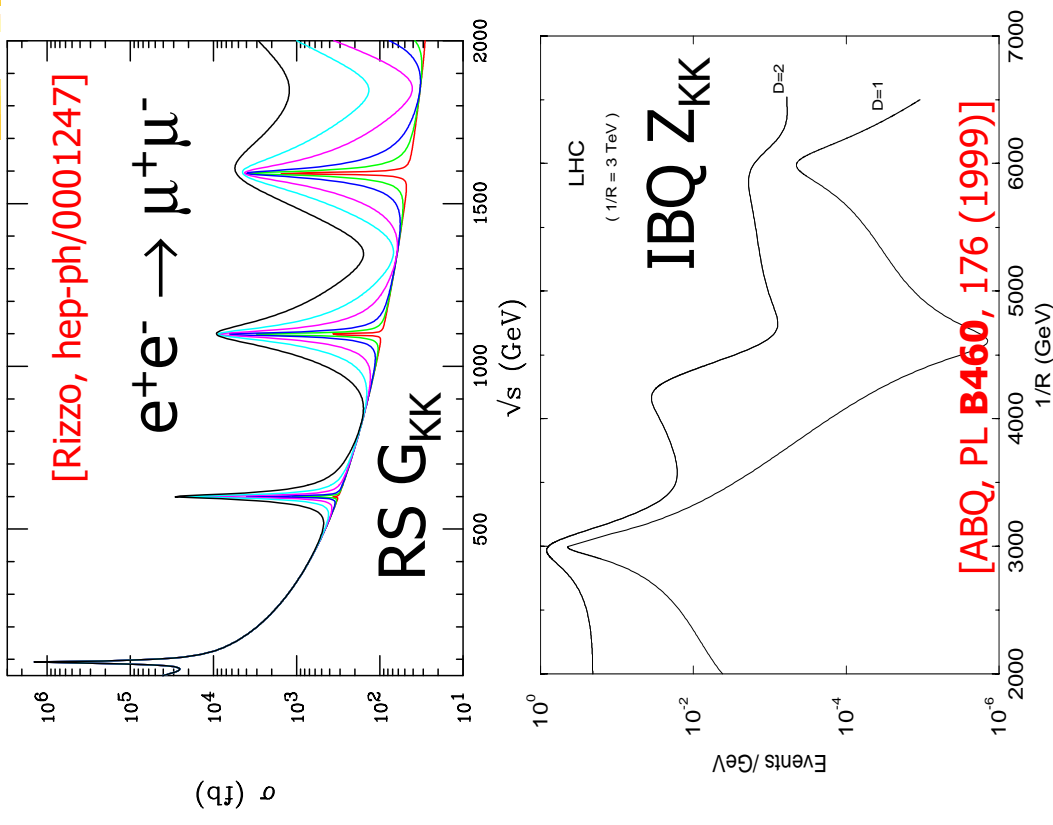
 New developments in theory of extra dimensions:

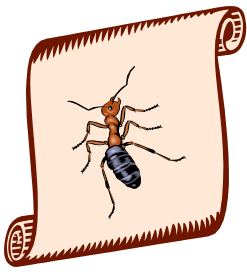
 **Randall-Sundrum** two-brane model with gravity localized near the brane [PRL **83**, 3370 (1999); PRL **83**, 4690 (1999)]

 Expect G_{KK} resonances in, e.g., $e^+e^- \rightarrow l^+l^-$ scattering

 **Antoniadis/Benaklis/Quiros** intermediate 'longitudinal' extra dimensions with $\sim \text{TeV}^{-1}$ radius [PL **B460**, 176 (1999)]

 Expect Z_{KK}, W_{KK}, g_{KK} resonances
 Effects also will be seen in the virtual exchange of the Kaluza-Klein modes of vector bosons at lower energies





Conclusion: WWW Search for Extra Dimensions

<http://www.extradimensions.com>

Coming Soon!

We recently registered our domain name...
register.cdn
 the first step on the job

Additional Services

- Domain Name Search
- URL Forwarding
- Web Services
- Free Domain Name Monitoring
- Sell Domain Name

CDW Computing Solutions
 Palm V Connected Organizer
 Long-life lithium-ion battery [click here](#)

Check for domain name availability:
 www.

Please be patient-it may take a minute.

© 1995-2000 by register.com™. Please read our [Disclaimer](#), [Registration Agreement](#), and [Dispute Policy](#).

Extra Dimensions TV Show

On 2/15/00 patent 6,025,810 was issued to David Strom for a "hyper-light-speed antenna." The concept is deceptively simple: "The present invention takes a transmission of energy, and instead of sending it through normal time and space, it pokes a small hole into another dimension, thus sending the energy through a place which allows transmission of energy to exceed the speed of light." According to the patent, this portal "allows energy from another dimension to accelerate plant growth." - from the AIP's "What's New", 3/17/00

- Stay tuned – next generation of collider experiments has a good chance to solve the mystery of large extra dimensions!
- Ultimate test of this exciting theory will become possible at the LHC