
Recoil and threshold resummation for hard scattering cross sections

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1 – Higher orders in hard scattering

Many hadronic sections in QCD factorize into parton distributions and hard scattering functions.

E.g. prompt photon production:

$$p_T^3 \frac{d\sigma_{AB \rightarrow \gamma X}(x_T^2)}{dp_T} = \sum_{ab} \phi_{a/A}(x_a, \mu) \otimes \phi_{b/B}(x_b, \mu) \otimes \omega_{ab \rightarrow \gamma X}(\hat{x}_T^2, \mu)$$

- $x_T^2 \equiv 4p_T^2/S$ and $\hat{x}_T^2 \equiv 4p_T^2/x_a x_b s$
- hard scattering function ω_{ab} computed in perturbation theory

Computation of higher order (α_s^k) corrections

- tests convergence of perturbation theory
- moves dependence on unphysical scales ($\mu_{F,R}$) to higher orders

Resummation of classes of log corrections

- may “rescue” perturbation theory if in distress
- moves dependence on unphysical scales to subleading terms
- may explain large data-theory discrepancies

(For good understanding, both should be done even if data don't seem to need it).

Two popular classes of large (Sudakov) logarithmic corrections (for Drell-Yan e.g.) $\frac{d\sigma}{dQ^2 dQ_T^2}$

1. Threshold enhancements

- $\alpha_s^i \left[\frac{\ln^{2i-1}(1-z)}{1-z} \right]_+$, $z = \frac{Q^2}{x_a x_b S}$ integration variable
- possibly large corrections after integrations against smooth functions (PDF's)

2. Q_T (recoil) enhancements

- $\alpha_s^i \frac{\ln^{2i-1}(Q_T/Q)}{Q_T}$, Q_T measurable
- large if $Q_T \ll Q$

Both result from simultaneously soft and collinear gluon radiation.

Can they be resummed together? (H-n.Li ('98))

2 – Refactorization and threshold resummation

Scheme (sch) dependence in Drell-Yan:

$$Q^2 \frac{d\sigma_{AB \rightarrow \gamma^* + X}}{dQ^2} = \sigma_0 \sum_{ab} \phi_{a/A}^{\text{sch}}(x_a, \mu) \otimes \phi_{b/B}^{\text{sch}}(x_b, \mu) \otimes \omega_{ab \rightarrow \gamma^* + X}^{\text{sch}} \left(z = Q^2 / (x_a x_b S), \frac{Q}{\mu} \right)$$

If we use $\overline{\text{MS}}$ -scheme PDF's:

$$\omega_{ab \rightarrow \gamma^* + X}^{(1), \overline{\text{MS}}} \simeq \alpha_s 4C_F \left[\frac{\ln(1-z)}{1-z} \right]_+$$

For DIS-scheme PDF's

$$\omega_{ab \rightarrow \gamma^* + X}^{(1), \text{DIS}} \simeq \alpha_s 2C_F \left[\frac{\ln(1-z)}{1-z} \right]_+$$

i.e. some logs \subset finite terms in the redefinition of the $\phi_{a/A}^{\text{sch}}$. Systematize [Sterman \('87\)](#) to factor **all** large corrections from ω_{ab} .

Define [Sterman\('87\)](#) the a' matrix element of **time-like** separated quark a fields.

$$\psi_{a/a'}(x, 2p_0, \epsilon) = \frac{1}{4\sqrt{2}N_c} \int \frac{d\lambda}{2\pi} e^{-i\lambda x p^0} \langle a'(p) | \bar{q}_a(\lambda, \vec{0}) \gamma^+ q_a(0) | a'(p) \rangle,$$

Fourier transform fixes **energy** fraction to be x .

$$\begin{aligned} \psi_{a/a'}(x, 2p_0, \epsilon) &= \delta(1-x) + \alpha_s \left(\frac{1}{\epsilon} P_{aa'}(x) \right. \\ &\quad \left. + 2C_F \left[\frac{\ln(1-z)}{1-z} \right]_+ + \dots \right) \end{aligned}$$

Define also U_{ab} : eikonalized ω_{ab} (with soft parts of ψ 's subtracted). Then refactorize hard scattering function

$$\begin{aligned} \omega_{ab \rightarrow \gamma^* + x} &= \psi_{a/a} \otimes \psi_{b/b} \otimes U_{ab} \\ &\quad \times H_{ab} + Y_{thr} \end{aligned}$$

H_{ab} has no large logs left! Large logs now in universal ψ 's. Also in U , but “easy”. “ \otimes ” includes energy conservation to $\mathcal{O}(1-z)$. Matching term $Y_{thr} = \mathcal{O}(1-z)$.

The ψ 's and U can be resummed to all orders, e.g.

$$\ln \psi(N) = \ln \int_0^1 dx x^{N-1} \psi(x) = \int_0^1 \frac{x^{N-1} - 1}{1-x} \left[\int_{(1-x)^2 Q^2}^{Q^2} \frac{d\mu^2}{\mu^2} A(\alpha_s(\mu)) + B((\alpha_s((1-x)Q)) \right]$$

Sterman('87); Catani, Trentadue('89)

- ψ 's universal, soft function U process dependent
- procedure extends to pure QCD cross sections, where color-coherence effects show up, Contopanagos, EL, Sterman ('97); Kidonakis, Oderda, Sterman ('98); Bonciani, Catani, Mangano, Nason ('98) and single particle inclusive kinematics EL, Oderda, Sterman('98)
- Accuracy achieved : NLL
- for $(1-x) < \Lambda/Q$ the μ integral diverges: "IR renormalon"

What are threshold-resummed cross sections good for?

- estimates of sizes of all-order cross sections (E.g. top quark cross section EL, Smith, van Neerven ('92,'94); Berger, Contopanagos ('98); Catani, Mangano, Nason, Trentadue ('96))
- generation of approximate higher orders (E.g. DY Magnea ('91); Higgs EL, Krämer, Spira ('98); F_2 A. Vogt ('99); $F_2(c\bar{c})$ EL, Moch ('99); $g_1(c\bar{c})$ Eynck, Moch ('00); prompt γ Kidonakis, Owens ('99); W, Z +jet Kidonakis, Del Duca ('99)).

All order cross sections need a prescription to deal with IR renormalon:

- cut-off Appel, Sterman, Mackenzie ('88): EL, Smith, van Neerven ('92)
- principal value Contopanagos, Sterman ('94)
- avoid in N contour integration Catani, Mangano, Nason, Trentadue ('96)

3 – Refactorization and double resummation

Write parton-level prompt photon p_T spectrum via distribution in (unobserved) Q_T of $a + b \rightarrow \gamma + c$ process [Lai, Li \('98\)](#)):

$$p_T^3 \frac{d\sigma_{ab \rightarrow \gamma c}^{(\text{resum})}(\bar{\mu})}{dp_T} \simeq \int d^2 \mathbf{Q}_T p_T^3 \frac{d\sigma_{ab \rightarrow \gamma c}^{(\text{resum})}}{d^2 \mathbf{Q}_T dp_T} \Theta(\bar{\mu} - Q_T) .$$

Refactorize

$$p_T^3 \frac{d\sigma_{ab \rightarrow \gamma c}^{(\text{resum})}}{d^2 \mathbf{Q}_T dp_T} = R_{a/a} \otimes R_{b/b} \otimes S_{ab} \otimes J_c \times h_{cd} + Y_{both}$$

with

$$R_{a/a}(x, \mathbf{k}, \epsilon) = \frac{1}{4\sqrt{2}N_c} \int \frac{d\lambda}{2\pi} \frac{d^2 b}{(2\pi)^2} e^{-i\lambda x p^0 + i\mathbf{b} \cdot \mathbf{k}} \langle a(p) | \bar{q}_a(\lambda, \mathbf{b}, 0) \gamma^+ q_a(0) | a(p) \rangle ,$$

Fourier transform fixes energy fraction *and* transverse momentum of quark a . “ \otimes ” includes both k_T and energy conservation. J_c is final state “jet” function (has only threshold logs).

J_c , S_{ab} and $R_{a/a}$ can be resummed (via eikonal approximation, and “webs” Gatherall ('83)):

$$\frac{p_T^3 d\sigma_{AB \rightarrow \gamma}^{(\text{resum})}}{dp_T} = \sum_{ij} \frac{p_T^4}{8\pi S^2} \int_c \frac{dN}{2\pi i} \tilde{\phi}_{i/A}(N, \mu) \tilde{\phi}_{j/B}(N, \mu) \int_0^1 d\tilde{x}_T^2 (\tilde{x}_T^2)^N \frac{|M_{ij}(\tilde{x}_T^2)|^2}{\sqrt{1 - \tilde{x}_T^2}} \times \int \frac{d^2 \mathbf{Q}_T}{(2\pi)^2} \Theta(\bar{\mu} - Q_T) \left(\frac{S}{4\mathbf{p}'_T{}^2} \right)^{N+1} P_{ij} \left(N, \mathbf{Q}_T, \frac{4p_T^2}{\tilde{x}_T^2}, \mu \right),$$

where

- $\mathbf{p}'_T = \mathbf{p}_T - \mathbf{Q}_T/2$
- $x_T^2 \equiv 4p_T^2/S$ and $\tilde{x}_T^2 \equiv 4\mathbf{p}_T^2/\hat{s}$
- “normal” PDF’s can be used!
- $P_{ij} \simeq \int d^2 b \exp(E_{ij}(N, \mathbf{b}, ..))$ “profile function”

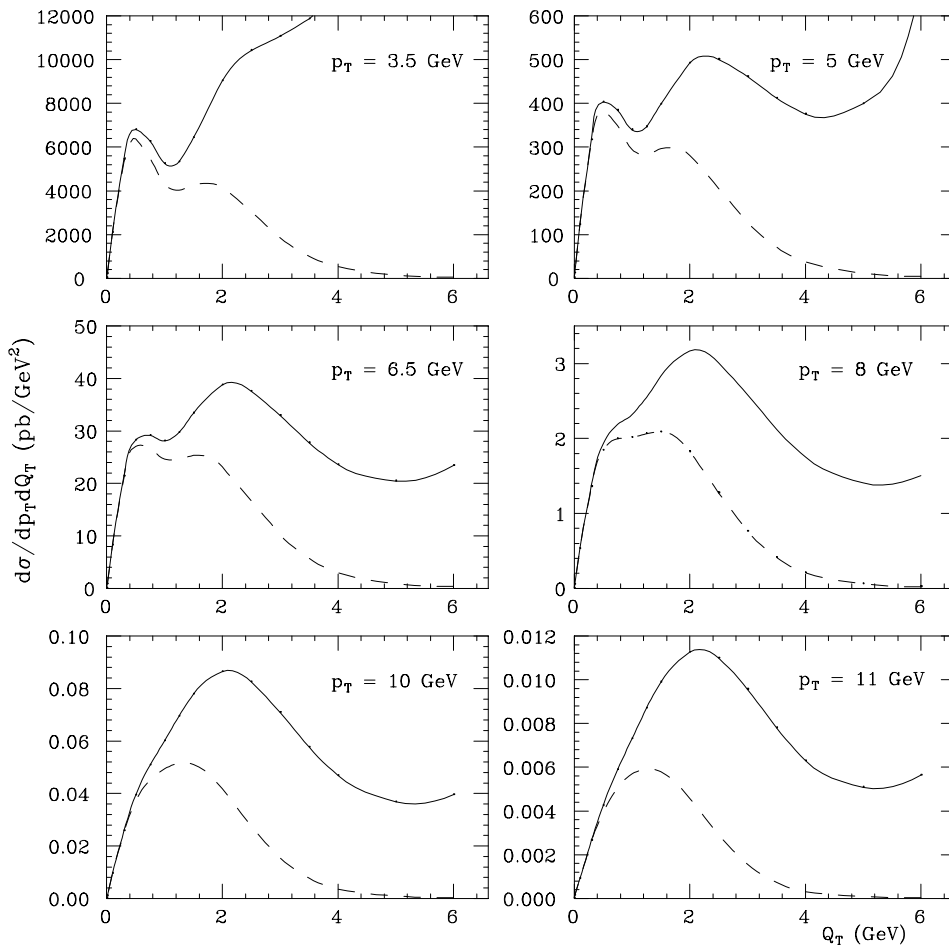
Here the “Drell-Yan $\ln(1 - z)$ ” is $\ln(1 - \tilde{x}_T^2)$.

What is double resummation good for?

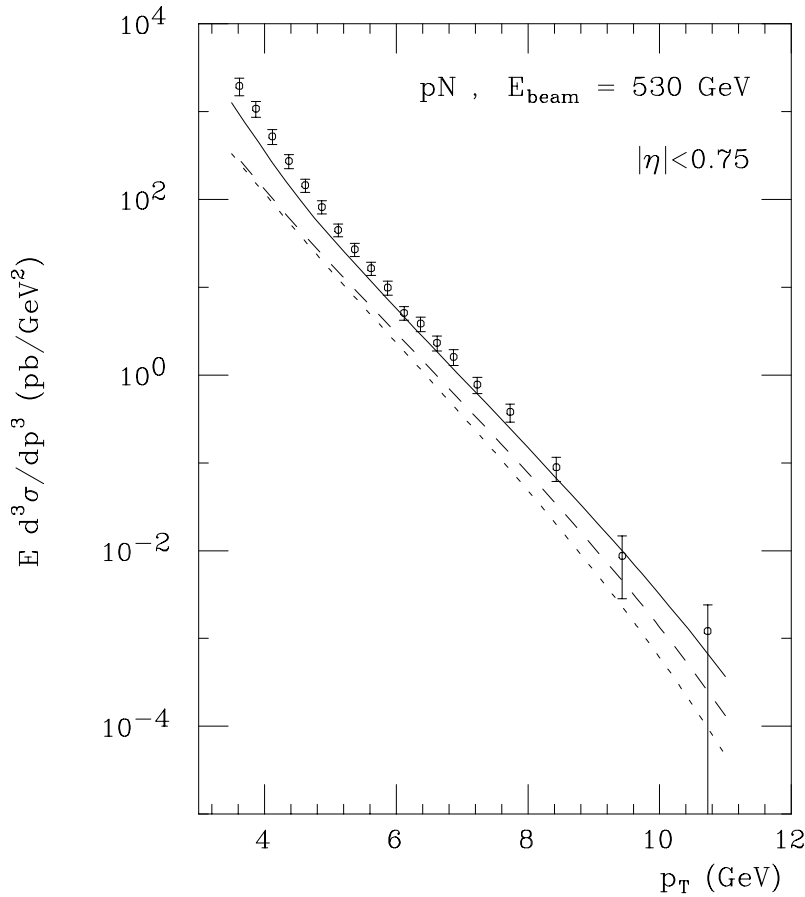
- estimates of shapes + sizes of all-order cross sections
- perturbative recoil effects included with energy conservation
- less need for phenomenological “intrinsic k_T ”?

4 – Is double resummation important?

Effect of profile function ($\sqrt{S} = 31.5$ GeV (E706)), vs. Q_T for $d\sigma_{pN \rightarrow \gamma X}^{(\text{resum})} / dQ_T dp_T$. Dashed lines: $(S/4p_T^2)^{N+1}$ (no recoil) instead of $(S/4\mathbf{p}'_T{}^2)^{N+1}$



Prompt photon cross section $E d^3 \sigma_{pN \rightarrow \gamma X} / dp^3$ for pN collisions at $\sqrt{s} = 31.5$ GeV. Data from E706, GRV PDF's.



- dotted: full NLO Aurenche et al ('88); Baer, Ohnemus, Owens ('90); Gordon, Vogelsang ('90)
- dashed: threshold resummation Catani, Mangano, Nason, Oleari, Vogelsang ('99); EL, Oderda, Sterman ('99)
- solid: double resummation

- Substantial numerical effect: for a measured p_T , the $2 \rightarrow 2$ scattering has to contribute less, with an accompanying less heavy final state.
- no matching yet. We took cutoff $\bar{\mu} = 5\text{GeV}$
- $\mu = p_T$, $F_{ij} = 0.5\text{GeV}^2$
- both resummed curves have substantially reduced μ dependence
- recoil of partons in hadrons evidently relevant!

5 – Summary and outlook

Summary:

- Recoil and threshold resummation now in one consistent formalism, to NLL accuracy
- formalism ready for Drell-Yan (Higgs..) at measured Q_T , and for single particle inclusive $2 \rightarrow 2$ scattering.
- (new principal value method for integration over impact parameter b)

Outlook:

- matching to finite order? Power corrections?
- other processes, numerical studies