

Constraints on R-parity
Violating Couplings from
Precision Electroweak
Measurements

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R-parity violation

MSSM

$$+ \frac{1}{2} \lambda_{ijk} L_i L_j E_k + \lambda'_{ijk} L_i Q_j \bar{D}_k$$

$$+ \frac{1}{2} \lambda''_{ijk} U_i \bar{D}_j \bar{D}_k$$

λ_{ijk} are already constrained to be small $O(10^{-2})$

\therefore violates lepton universality in lepton decays at tree level.

Try to constrain λ'_{ijk} and λ''_{ijk} .

⇒ Look at Z-pole or W-pole observables only.

- ★ Sensitive to Vertex Corrections
- ★ Blind to Oblique Corrections
- ★ No condition on the new physics scale.
- ★ Can accommodate mixing with extra gauge bosons.
(treat as vertex correction)
- ★ No need to specify the model completely
- ★ Can lead to strong constraints

Z pole observables :

10
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m_Z , ~~Γ_Z~~

$$\sigma_{had}^0 = \frac{12\pi}{m_Z^2} \frac{\Gamma_{e^+e^-} \Gamma_{had}}{\Gamma_Z^2}$$

$$R_e = \frac{\Gamma_{had}}{\Gamma_{e^+e^-}} \quad (l = e, \mu, \tau)$$

$$A_f = \frac{(g_L^+)^2 - (g_R^+)^2}{(g_L^+)^2 + (g_R^+)^2} \quad (f = e, \mu, \tau, s, c, b)$$

$$A_{FB}(f) = \frac{3}{4} A_e A_f \quad (f = e, \mu, \tau, u, s, c, b)$$

$$R_b = \frac{\Gamma_b}{\Gamma_{had}}, \quad R_c = \frac{\Gamma_c}{\Gamma_{had}}$$

$$R_s' = \frac{\Gamma_s}{\Gamma_u + \Gamma_d + \Gamma_s}$$

Ratios of coupling constants

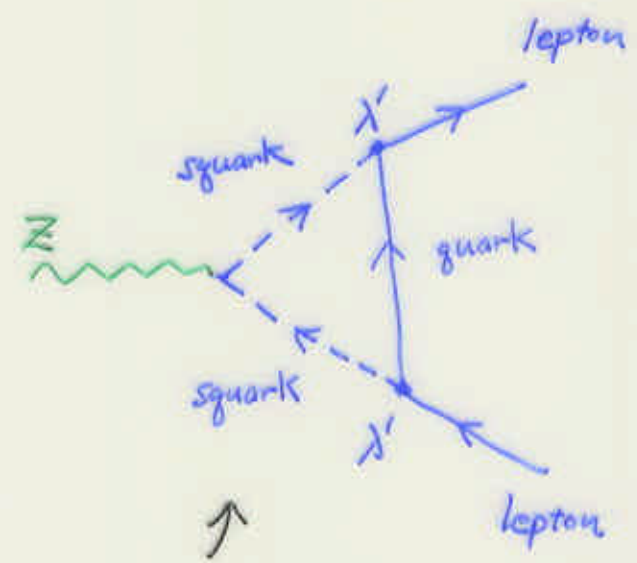
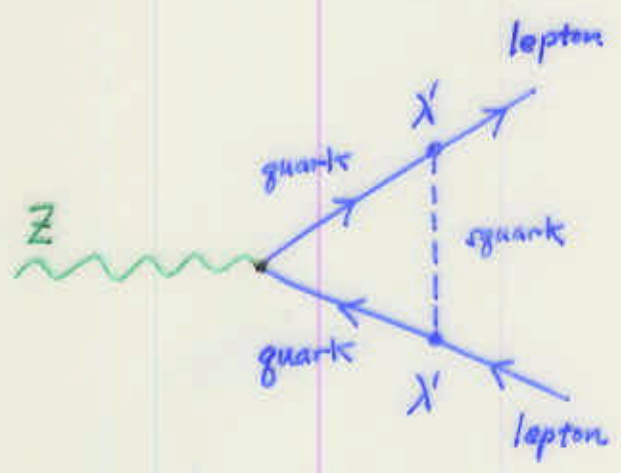
⇒ Depends on Oblique Corrections
only through $\sin^2 \theta_w$

Simplifying Approximations :

- ★ Take all quarks and leptons to be massless except the top.
 - ★ Take all slepton masses to be degenerate.
 - ★ Take all squark masses to be degenerate.
- ⇒ flavor dependence ^{of W and Z couplings} only from
- R-parity violating couplings and
- Higgs coupling to the left-handed bottom.

$$\delta_{ijk} L_i Q_j D_k$$

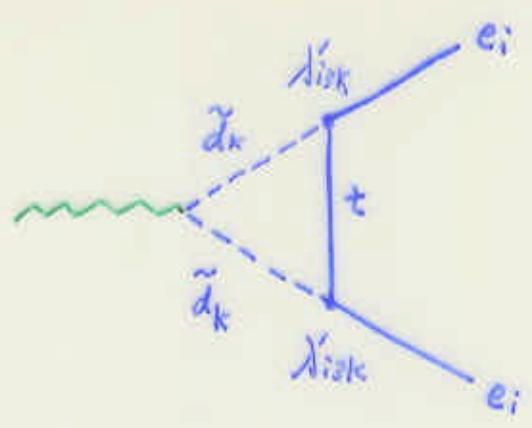
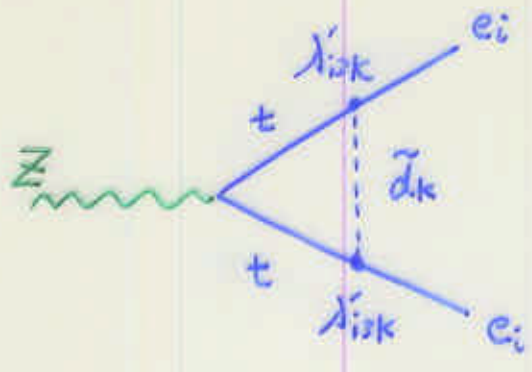
↑ corrects W and Z coupling to i-th generation left handed lepton.



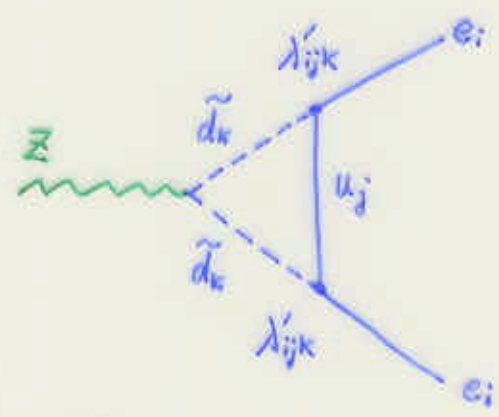
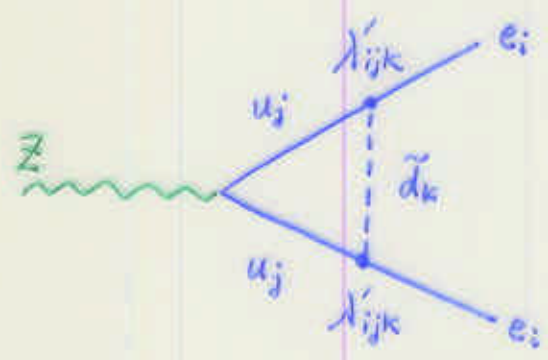
↑ depend on quark and squark masses

↓

Must consider the diagrams with the top ($j=3$) separately.

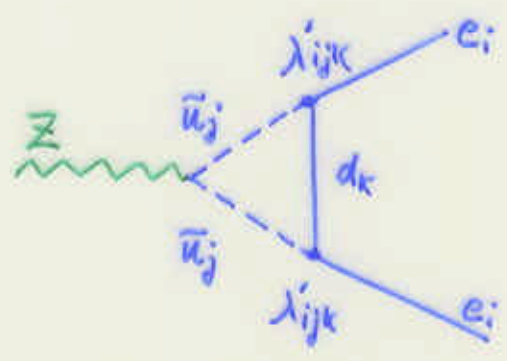
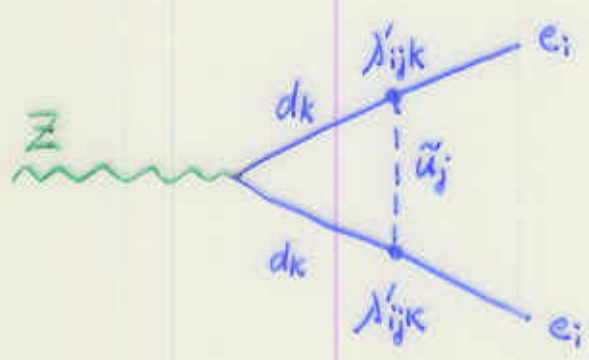


$$\delta h_{i3k}^{(t)} = + 0.63\% |\lambda'_{i3k}|^2 \quad \leftarrow \text{Square mass} = 100 \text{ GeV}$$



(j ≠ 3)

$$\delta h_{ijk}^{(u)} = - 0.02\% |\lambda'_{ijk}|^2$$



$$\delta h_{ijk}^{(d)} = - 0.06\% |\lambda'_{ijk}|^2$$

★ Keep only top-quark diagrams.

15
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$$\delta h_i^{\uparrow} \approx \sum_K \delta h_{i\uparrow K}^{(u)} = +0.63\% \sum_K |\lambda'_{i\uparrow K}|^2$$

↑
Shift in left handed coupling of i-th lepton

$$R_l = \frac{T_{had}}{T_{le^-}} = \frac{N_c \sum_f (h_{gl}^2 + h_{gr}^2)}{h_{ll}^2 + h_{lr}^2} \quad (l=e, \mu, \tau)$$

$$\frac{\delta R_l}{R_l} = \Delta_R + 4.3 \delta h_l^{\uparrow}$$

$$A_l = \frac{h_{al}^2 - h_{ar}^2}{h_{al}^2 + h_{ar}^2}, \quad A_{FB}(l) = \frac{3}{4} A_e A_l$$

$$\frac{\delta A_l}{A_l} = \Delta_A - 25 \delta h_l^{\uparrow}$$

$$\frac{\delta A_{FB}(l)}{A_{FB}(l)} = 2\Delta_A - 25 \delta h_e^{\uparrow} - 25 \delta h_l^{\uparrow}$$

5 parameters : $\Delta_R, \Delta_A, \delta h_e^{\uparrow}, \delta h_\mu^{\uparrow}, \delta h_\tau^{\uparrow}$

$$\left\{ \begin{array}{l} \Delta_{Re} \equiv \Delta_R + 4.3 \delta h_e^{\pi} \\ \Delta_{Ae} \equiv \Delta_A - 25 \delta h_e^{\pi} \\ \delta_{\mu e} \equiv \delta h_{\mu}^{\pi} - \delta h_e^{\pi} \\ \delta_{\tau e} \equiv \delta h_{\tau}^{\pi} - \delta h_e^{\pi} \end{array} \right.$$

↓

$$\left\{ \begin{array}{l} \Delta_{Re} = 0.0007 \pm 0.0020 \\ \Delta_{Ae} = 0.052 \pm 0.012 \\ \delta_{\mu e} = 0.00038 \pm 0.00056 \\ \delta_{\tau e} = -0.00013 \pm 0.00061 \end{array} \right.$$

↓

$$\left\{ \begin{array}{l} \sum_k |\lambda'_{23k}|^2 - \sum_k |\lambda'_{13k}|^2 = 0.062 \pm 0.095 \\ \sum_k |\lambda'_{33k}|^2 - \sum_k |\lambda'_{13k}|^2 = -0.02 \pm 0.10 \end{array} \right.$$

Neglect λ'_{12K} (already well constrained)

①⑦
②

$$\begin{cases} \sum_K |\lambda'_{23K}|^2 \leq 0.16 \quad (0.25) \\ \sum_K |\lambda'_{33K}|^2 \leq 0.08 \quad (0.18) \end{cases}$$

↓

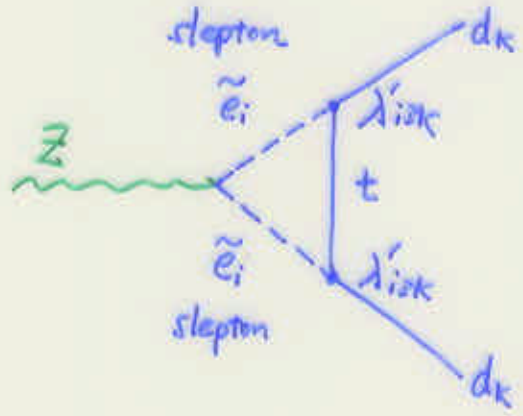
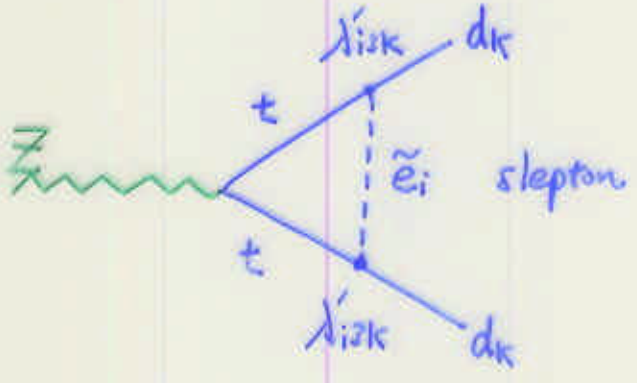
$$\begin{cases} |\lambda'_{23K}| \leq 0.40 \quad (0.50) \\ |\lambda'_{33K}| \leq 0.28 \quad (0.42) \end{cases}$$

Limit from W -decay

$$|\lambda'_{32K}| \leq 1.7 \quad (2.4)$$

$$\lambda'_{ijk} L_i Q_j D_k$$

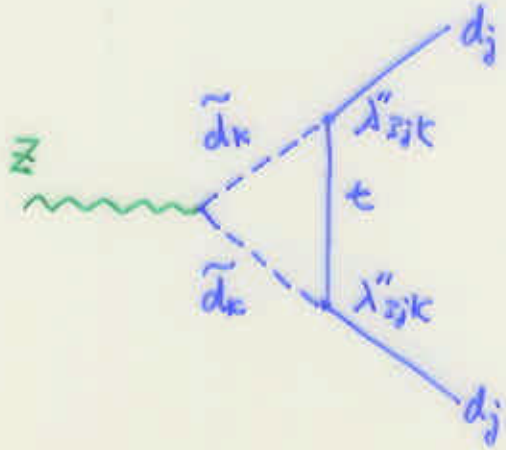
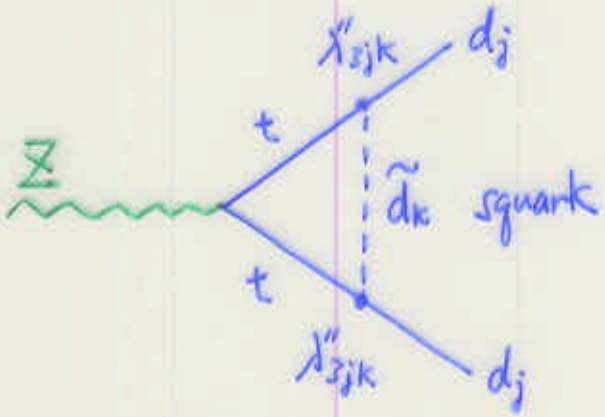
only top \uparrow \nwarrow corrects right-handed down-type quark couplings.



$$\frac{1}{2} \lambda''_{ijk} U_i D_j D_k$$

top \uparrow antisymmetric \nwarrow

corrects right-handed down-type quark couplings



Flavor Dependent shifts :

$$\delta h_d^{\pi}, \delta h_s^{\pi}, \delta h_u^{\pi}, \delta h_b^{\text{Higgs}}$$

↑
shift to left-handed
coupling of the b.

Flavor Independent shifts :

$$\delta s^2$$

↑ shift in effective value of $\sin^2 \theta_w$.

5 parameter fit

$$\left\{ \begin{array}{l} \delta s^2 = -0.00092 \pm 0.00022 \\ \delta h_d^{\pi} = 0.081 \pm 0.077 \\ \delta h_s^{\pi} = 0.055 \pm 0.043 \\ \delta h_u^{\pi} = 0.026 \pm 0.010 \\ \delta h_b^{\text{Higgs}} = -0.0031 \pm 0.0042 \end{array} \right.$$

$$\delta h_{dk}^k = \underline{-0.215\%} \sum_i |\lambda'_{i3k}|^2$$

$$= \underline{-0.43\%} \sum_j |\lambda''_{j3k}|^2$$

$$\left\{ \begin{array}{l} \sum_i |\lambda'_{i31}|^2 \leq -2 \quad (34) [70] \\ \sum_i |\lambda'_{i32}|^2 \leq -6 \quad (14) [34] \\ \sum_i |\lambda'_{i33}|^2 \leq -7.4 \quad (-2.8) [1.9] \end{array} \right. \rightarrow \left\{ \begin{array}{l} |\lambda'_{i31}| \leq (5.8) [8.4] \\ |\lambda'_{i32}| \leq (3.8) [5.9] \\ |\lambda'_{i33}| \leq (-) [1.4] \end{array} \right.$$

$$\left\{ \begin{array}{l} \sum_k |\lambda''_{31k}|^2 \leq -1 \quad (17) [35] \\ \sum_k |\lambda''_{32k}|^2 \leq -3 \quad (7) [17] \\ \sum_k |\lambda''_{33k}|^2 \leq -3.7 \quad (-1.4) [0.9] \end{array} \right. \rightarrow \left\{ \begin{array}{l} |\lambda''_{321}| \leq (2.7) [4.1] \\ |\lambda''_{33i}| \leq (-) [0.96] \end{array} \right.$$

Conclusions

- ☆ To go beyond the standard STU analysis, look at vertex corrections at the Z and W poles.
- ☆ Application of the technique places strong constraints on R -parity violating couplings.