

THE

STANDARD MODEL

PREDICTION

FOR

ϵ'/ϵ

OUTLINE

- FSI IN KAON DECAYS
 - CONCEPTS
 - THE OMNES PROBLEM
 - PION SCALAR FORM FACTOR
 - $K \rightarrow \pi\pi$ AMPLITUDES

- SM ESTIMATE OF ϵ'/ϵ

REFS. :

A. PICH, E.P. hep-ph/9911233 PRL 84 (2000)

E.P. hep-ph/0007017

A. PICH, E.P. hep-ph/0007208

A. PICH, I. SCINEMI, E.P. to appear

The present experimental world average is

$$\text{Re}(\varepsilon'/\varepsilon) = (19.3 \pm 2.4) \times 10^{-4}$$

SM predictions are systematically below

$$(\varepsilon'/\varepsilon)_{\text{SM}} \approx 7 \times 10^{-4} \pm ?$$

[S. Bosch et al.
hep-ph/9904408]

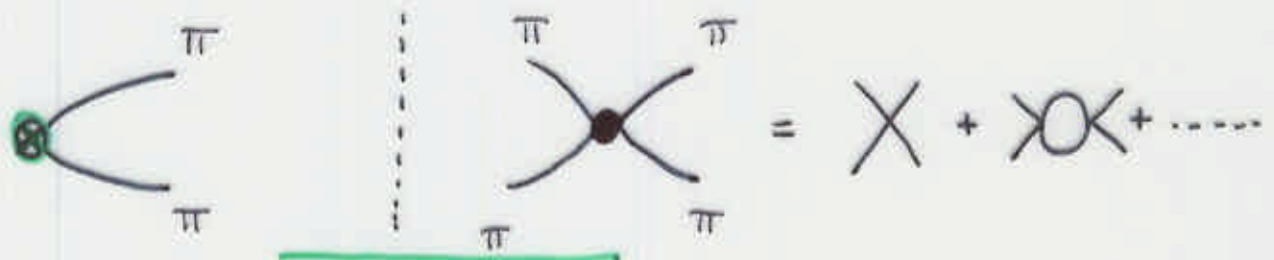
But FSI effects are not included

$K \rightarrow \pi\pi$ at one-loop in CHPT gives



FSI are large for $I=0$

FSI - CONCEPTS



$$J^P = 0^+ \quad I = 0$$

PRODUCTION

RESCATTERING

- universal
- no factorization scale dependence
- well known

$$A_I \equiv A[k \rightarrow (\pi\pi)_I] = A_I e^{i\delta_0^I} \quad I = 0, 2$$

$$(\delta_0^0 - \delta_0^2)(m_K^2) = 45^\circ \pm 6^\circ \quad [\text{GASSER, HEISSNER 91}]$$



ANALITICITY + UNITARITY

THE OHNÈS PROBLEM

$$A(s) = \langle (\pi\pi)_e^I | 0 | i \rangle \quad s = P_e^2$$

$$L = [4m_\pi^2, \infty)$$

In the Elastic region + Watson's Theorem

$$A(s) = |A(s)| e^{i\delta_e^I(s)}$$

$$\text{Im } A(s) = \delta^T e^I A^* = e^{i\delta_e^I} \sin \delta_e^I A^*$$

Dispersion relation

$$A(s) = \frac{1}{\pi} \int_L dz \frac{\text{Im } A(z)}{z - s - i\varepsilon} + \text{SUBTRACTIONS}$$

Omnès Solution [s_0 = subtraction point]

$$A(s) = Q_M(s, s_0) \Omega_M(s, s_0)$$

↪ OMNÈS FACTOR

$$\begin{aligned} \Omega_M(s, s_0) &= \exp \left\{ \frac{(s-s_0)^M}{\pi} \int_{\mathcal{L}} \frac{dz}{(z-s_0)^M} \frac{\delta_e^{\mathbb{I}}(z)}{z-s-i\varepsilon} \right\} \\ &\equiv e^{i\delta_e^{\mathbb{I}}(s)} \cdot \mathcal{R}_M(s, s_0) \end{aligned}$$

v.B. $\Omega_M(s_0, s_0) = 1$

$$\ln Q_M(s, s_0) = \sum_{k=0}^{M-1} \frac{(s-s_0)^k}{k!} \left. \frac{\partial^{(k)} \ln A}{\partial s^{(k)}} \right|_{s=s_0}$$

$F_S^\pi(t)$: SCALAR FORM FACTOR

$$\langle \pi^i | \bar{u}u + \bar{d}d | \pi^k \rangle = \delta^{ik} F_S^\pi(t)$$

$$\boxed{I=0}$$

Known at one loop (up to two loops) in ChPT

$$F_S^\pi(t) = F_S^\pi(0) \{ 1 + g(t) + O(p^4) \}$$

The Omnès solution gives:

$$F_S^\pi(t) = \Omega_0(t, t_0) \cdot \bar{F}_S^\pi(t_0) \approx \Omega_0(t, t_0) \cdot F_S^\pi(0) [1 + g(t_0)]$$

where $g(t_0) \rightarrow 0$ for $t_0 \rightarrow 0$

$$\Omega_0(t, t_0) \equiv e^{i\delta_0^0(t)} \cdot \mathcal{R}_0(t, t_0)$$

A twice-subtracted analysis gives

$$\mathcal{R}_0(M_K^2, 0) = 1.55 \pm 0.10$$

It takes into account:

- Inelastic contributions ($t > 1 \text{ GeV}^2$)
- Experimental phase-shifts uncertainties

K \rightarrow $\pi\pi$

- The Omnès solution

$$\boxed{I = 0, 2}$$

$$A_I(M_K^2) = (M_K^2 - M_\pi^2) \cdot \Omega_I(M_K^2, s_0) \cdot a_I(s_0)$$

where $a_I(s) = a_I(0) \left\{ 1 + g_I(s) + O(p^4) \right\}$

↑
dominated by
FSI (IR) effects

- The Omnès factors at $s_0 = 0$

$$\boxed{R_0(M_K^2, 0) = 1.55 \pm 0.10}$$

$$\boxed{R_2(M_K^2, 0) = 0.92 \pm 0.03}$$

$$\begin{array}{ccc} \uparrow & \uparrow & \uparrow \\ \text{I}=0 & \text{I}=2 & \text{I}=2 \end{array}$$

("LARGE N_c ")
MATCHING
EXACT

$$\times \mathcal{R}_0 \quad \times \mathcal{R}_2$$

$$\times \mathcal{R}_2$$

The SM prediction is enhanced to

$$\text{Re}(\varepsilon'/\varepsilon)_{\text{SM}} = (17 \pm 6) \times 10^{-4}$$

ϵ'/ϵ

EXP
↓

THEORY
↓

EXP.
↓

THEORY
↓

$$\epsilon'/\epsilon \approx \tau \sum_i y_i(\mu) \langle Q_i(\mu) \rangle_0 (1 - \Omega_{10}) - \frac{\tau}{\omega} \sum_i y_i(\mu) \langle Q_i(\mu) \rangle_2$$

$$N_c \rightarrow \infty$$

$$d_s \sim \frac{1}{N_c}$$

$$\ln \frac{m_c}{\mu} \text{ SMALL}$$



$$\langle Q(\mu) \rangle_{\mathbf{I}} \sim (\text{LARGE } N_c) \times R_{FSI}^{\mathbf{I}}$$

ENERGY SCALE

FIELDS

EFFECTIVE THEORY

